LINEAR PROGRAMMING BOUNDS FOR HYPERBOLIC SURFACES

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ABSTRACT. We adapt linear programming methods from sphere packings to closed hyperbolic surfaces and obtain new upper bounds on their systole, their kissing number, the first positive eigenvalue of their Laplacian, the multiplicity of their first eigenvalue, and their number of small eigenvalues. Apart from a few exceptions, the resulting bounds are the current best known both in low genus and as the genus tends to infinity. Our methods also provide lower bounds on the systole (achieved in genus 2 to 7, 14, and 17) that are sufficient for surfaces to have a spectral gap larger than 1/4.



FIGURE 1. Upper bounds and current record holders for the maximization of geometric and spectral invariants associated to hyperbolic surfaces.

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1. INTRODUCTION

The goal of this paper is to prove new upper bounds on five invariants associated to a closed, oriented, hyperbolic surface M:

- (1) its systole sys(M), the length of any shortest non-contractible closed curve in M;
- (2) its kissing number kiss(M), the number of homotopy classes of oriented noncontractible closed curves of minimal length in M;
- (3) the first positive eigenvalue $\lambda_1(M)$ (which coincides with the spectral gap) of the Laplace-Beltrami operator Δ_M on M;
- (4) the multiplicity $m_1(M)$ of the eigenvalue $\lambda_1(M)$, that is, the dimension of the corresponding eigenspace;
- (5) the number $N_{\text{small}}(M)$, counting multiplicity, of small eigenvalues of Δ_M , that is, those contained in the interval [0, 1/4].

We bound the first four invariants in terms of the genus of M only, but the fifth one in terms of the genus and the systole. In low genus (or for small systole), our bounds are illustrated in Figures 1 and 2 (see also Tables 1 to 7). They beat all previous upper bounds except for the systole and kissing number in genus 2, for λ_1 in genus 2, 3, 4, and 6, and for N_{small} when the systole is smaller than 2.317. Note that $N_{\text{small}}(M) < 2$ if and only if $\lambda_1(M) > 1/4$. We use this to show that there exist surfaces with a spectral gap larger than 1/4 in genus 4 to 7, 14, and 17 (this was already known in genus 2 and 3). Whether such surfaces exist in every genus is a well-known open problem related to Selberg's eigenvalue conjecture [Sel65] (see e.g. [Mon15, Question 1.1] and [Wri20, Problem 10.4]). A lot of progress on this question in high genus was made recently [MNP22, WX22, LW21, HM21].

In higher genus (or for larger systole), our asymptotic bounds are as follows.

Theorem 6.4. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$sys(M) < 2\log(g) + 2.409.$$

Theorem 7.8. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

kiss
$$(M) < \frac{4.873 \cdot g^2}{\log(g) + 1.2045}$$

Theorem 8.3. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$\lambda_1(M) < \frac{1}{4} + \left(\frac{\pi}{\log(g) + 0.7436}\right)^2$$

Theorem 9.5. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$m_1(M) \le 2g - 1.$$

Theorem 10.2. If M is a closed hyperbolic surface of genus $g \ge 2$, then

$$N_{\text{small}}(M) < \min\left(\frac{24\pi^2(g-1)}{\text{sys}(M)^3}, \frac{16(g-1)}{\text{sys}(M)^2}\right).$$

These improve upon the previous best upper bounds established in [Bav96], [FBP22] (previously [Par13]), [Che75], [Sév02], and [Hub76] respectively. While the previous bounds used very different techniques from one invariant to another, our proofs are all based on the same method, namely, linear programming.



FIGURE 2. Upper bounds on the number of small eigenvalues and lower bounds on the systole that imply a spectral gap larger than a quarter.

In addition to finding new examples of surfaces with $\lambda_1 > 1/4$, we also improve the previous best lower bounds of Colbois and Colin de Verdière [CCdV88] on the maximum of m_1 in genus 4, 7, 8, 10, 14 to 16, and 19.

Context. The invariants sys, kiss, λ_1 , m_1 , and N_{small} can be defined for any closed Riemannian manifold (with 1/4 replaced by the bottom of the spectrum of the Laplacian on the universal cover of M) and their maximization has been studied by several authors for various classes of manifolds. For example, if we fix a topological manifold Σ , then maximizing the systole over all Riemannian manifolds M of a given volume that are homeomorphic to Σ is called the *isosystolic problem* and it has been solved for the projective plane [Pu52], the 2-dimensional torus (Loewner), and the Klein bottle [Bav86]. The maximum of m_1 is known for the same surfaces [Bes80, CdV87, Nad88] as well as for the 2-sphere [Che75], while the maximum of λ_1 is known for the closed orientable surface of genus 2 [JLN⁺05, NS19] in addition to all the previous surfaces [Her70, LY82, Nad96, ESGJ06].

Another much-studied case is that of flat d-dimensional tori of unit volume. In that case, $\lambda_1(M) = (2\pi \operatorname{sys}(M^*))^2$ and $m_1(M) = \operatorname{kiss}(M^*)/2$ where M^* is the torus dual to M, so the first four maximization problems above reduce to only two while the fifth is trivial since the bottom of the spectrum of the Laplacian on \mathbb{R}^d is 0. Furthermore, if Λ is a lattice such that $M \cong \mathbb{R}^d / \Lambda$, then the balls of radius $\operatorname{sys}(M)/2$ centered at the points in Λ form a sphere packing of density $2^{-d} \operatorname{sys}(M)^d \operatorname{vol}(B_1^d)$, where B_1^d is the unit ball in \mathbb{R}^d , and there are exactly $\operatorname{kiss}(M)$ balls tangent to any ball in this packing. In other words, maximizing the systole and kissing number of flat tori of unit area is equivalent to maximizing the packing density and kissing number among sphere packings whose centers form a lattice. Both problems have been solved in dimensions 1 to 8 and 24 (see [CS93, Table 1.1] and [CK09]). In some cases, the solutions to these problems were obtained by a

method known as *linear programming* first introduced by Delsarte in the context of errorcorrecting codes [Del72]. This was then adapted to prove bounds on the kissing numbers of arbitrary sphere packings (not necessarily coming from lattices) in [DGS77] and then on their packing density in [CE03]. In addition to giving optimal bounds in dimensions 8 [Via17] and 24 [CKM⁺17], this approach also yields the best known asymptotic bounds on kissing numbers and packing density as the dimension tends to infinity [CZ14].

The five invariants we consider here were previously investigated for hyperbolic surfaces in [Hub74, Che75, Hub76, Bus77, Hub80, Jen84, BC85, Bro88, BBD88, CCdV88, Bur90, Sch93, Sch94, BS94, Bav96, Bav97, SS97, Ada98, Ham01, HK02, Kim03, CB05, KSV07, Ota08, Gen09, OR09, Par13, SU13, FP15, Gen15, Coo18, PW18, Pet18, HM21, Jam21, KMP21, LW21, Bon22, FBR22, Mon22, MNP22, WX22] among many others. For closed surfaces, the only optimal bounds known to date are for the systole [Jen84] and kissing number [Sch94] in genus 2, for m_1 in genus 3 [FBP21], and for $N_{\rm small}$ in every genus [OR09]. However, the known examples that maximize $N_{\rm small}$ have a short pants decomposition [Bus77], which is why we are interested in improved bounds as the systole grows.

We previously adapted the linear programming method to hyperbolic surfaces in [FBP22] and [FBP21] to prove bounds on kissing numbers in large genus and on m_1 in small genus respectively. Here we extend and improve our previous approach in a systematic way.

Organization. The paper is organized as follows. We start with preliminary sections on the Fourier transform, the Selberg trace formula, the linear programming method, and Bessel functions. This is followed by one section for each of the five invariants sys, kiss, λ_1 , m_1 , and N_{small} . In each of these sections, we first present a general criterion for proving upper bounds based on the Selberg trace formula. We then discuss the results that we have obtained from this criterion in low genus (or for small systole) using numerical optimization and conclude each section by proving an asymptotic bound. In both ranges, we compare our bounds with the previous best. The lower bounds on the systole that are sufficient to obtain a spectral gap larger than a quarter are described in subsection 10.4 and the new examples with large m_1 are presented in subsection 9.2.

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2. The Fourier transform

The Fourier transform of an integrable function $f: \mathbb{R} \to \mathbb{C}$ is defined by

$$\widehat{f}(y) := \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x) e^{-iyx} dx$$

for $y \in \mathbb{R}$. If \hat{f} is integrable, then the Fourier inversion theorem says that its Fourier transform is almost everywhere equal to $x \mapsto f(-x)$.

We will frequently use the scaling property that for a > 0, the Fourier transform of $x \mapsto f(ax)$ is $y \mapsto \widehat{f}(y/a)/a$. If f and g are integrable, then

$$\widehat{f \ast g} = \sqrt{2\pi} \,\widehat{f} \,\widehat{g}$$

where * denotes the convolution and if \hat{f} and \hat{g} are also integrable, then

$$\widehat{fg} = \frac{1}{\sqrt{2\pi}}\widehat{f} * \widehat{g}.$$

When f is an even function, which will be the case throughout the paper, its Fourier transform reduces to the cosine transform

$$\widehat{f}(y) = \sqrt{\frac{2}{\pi}} \int_0^\infty f(x) \cos(yx) \, dx$$

and is therefore even.

An integrable function f with integrable Fourier transform is said to be *positive-definite* if $\hat{f}(y) \geq 0$ for every $y \in \mathbb{R}$. This is not the usual definition, but is equivalent to it by Bochner's theorem. The set of positive-definite functions is closed under convolution and multiplication.

3. The Selberg trace formula

Given a closed hyperbolic surface M (always assumed to be oriented), we list the eigenvalues of its Laplace–Beltrami operator (acting on square-integrable functions) in nondecreasing order

$$0 = \lambda_0(M) < \lambda_1(M) \le \lambda_2(M) \le \cdots$$

where each eigenvalue is repeated according to its multiplicity, i.e., the dimension of the corresponding eigenspace.

The set of oriented closed geodesics in M is denoted by $\mathcal{C}(M)$. This means that each unoriented closed geodesic appears twice in $\mathcal{C}(M)$, once for each orientation. The length of a geodesic $\gamma \in \mathcal{C}(M)$ is denoted $\ell(\gamma)$ and its *primitive length* is denoted $\Lambda(\gamma)$. The latter is defined as the length of the shortest geodesic α such that $\gamma = \alpha^k$ for some power $k \ge 1$. The geodesic α is called *primitive* because it cannot be expressed as a proper power of another geodesic.

A function $f : \mathbb{R} \to \mathbb{C}$ is said to be *admissible* if it is even, integrable, and its Fourier transform \hat{f} is holomorphic in a horizontal strip of the form

$$\left\{z \in \mathbb{C} : |\operatorname{Im}(z)| < \frac{1}{2} + \varepsilon\right\}$$

for some $\varepsilon > 0$ and satisfies the decay condition

$$|\widehat{f}(z)| = O\left(\frac{1}{1+|z|^p}\right)$$

for some p > 2 in that strip. Note that the decay condition implies that \hat{f} and $y\hat{f}(y)$ are integrable on the real line so that f itself must be continuously differentiable by Fourier inversion and differentiation under the integral sign.

With the above notation and normalizations, the Selberg trace formula [Bus10, Section 9.5] states that for every closed hyperbolic surface M of genus g and every admissible function $f : \mathbb{R} \to \mathbb{C}$, we have

$$\sum_{j=0}^{\infty} \widehat{f}\left(\sqrt{\lambda_j(M) - \frac{1}{4}}\right) = 2(g-1) \int_0^{\infty} \widehat{f}(x) x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(M)} \frac{\Lambda(\gamma) f(\ell(\gamma))}{2\sinh(\ell(\gamma)/2)}$$

Since \hat{f} is even, it does not matter which square root we use on the left-hand side. It is customary to write $r_j(M)$ for either of the two roots, so that $r_j(M)^2 + \frac{1}{4} = \lambda_j(M)$. Note that our convention for the Fourier transform differs from the one used in [Bus10] by a factor of $1/\sqrt{2\pi}$, which explains the appearance of this factor in the above formula.

4. LINEAR PROGRAMMING

Like the linear programming bounds of Cohn and Elkies [CE03] for the density of sphere packings, for each of the five invariants sys, kiss, λ_1 , m_1 , and N_{small} , our criterion will take the following form:

Suppose that f is an admissible function such that f and \hat{f} satisfy certain linear inequalities over certain intervals. Then f and \hat{f} produce a bound on the given invariant that holds for every closed hyperbolic surface (satisfying certain conditions) in a given genus.

This is called "linear programming" because the inequalities are linear in the sense that any positive linear combination of functions that satisfy the inequalities still satisfies the inequalities. However, the function to be optimized (the resulting bound) is not linear in f. Moreover, the space we are optimizing over is infinite-dimensional and there are infinitely many inequalities to check (one at each point in the specified intervals). For these reasons, classical linear programming algorithms do not work well, which led Cohn and Elkies to devise the following strategy (adapted here to our setting).

The idea is to consider functions f of the form $f(x) = p(x^2)e^{-x^2/2}$ where p is a polynomial. Such a function is automatically admissible since its Fourier transform takes the form $\hat{f}(y) = q(y^2)e^{-y^2/2}$ for some polynomial q, hence defines an entire function with super-exponential decay in any horizontal strip. Moreover, the map $p \mapsto q$ is linear. In fact, it is diagonal with entries $(-1)^n$ with respect to the basis of generalized Laguerre polynomials $L_n^{(-1/2)}$. The upshot is that it is possible to impose linear equations on both f and \hat{f} simultaneously. All one has to do is solve a linear system of equations to find the coefficients of p and q. The conditions we impose are that f and \hat{f} have double zeros at certain points x_1, \ldots, x_m and y_1, \ldots, y_n respectively.

The reason for imposing double zeros is that it prevents local changes of sign and with enough double zeros at appropriate locations we are usually able to find some functions fand \hat{f} that satisfy the required inequalities. Once we find such suitable zeros we then try to wiggle them to decrease the resulting bound, then add more zeros and repeat.

For sphere packing bounds this scheme appears to converge quickly to a unique optimal function f in each dimension. We have not found this to be the case for hyperbolic surfaces. One important difference is that for sphere packings, Cohn and Elkies assume that f and \hat{f} have the same double zeros and the situation is fairly symmetric. This is not the case with the Selberg trace formula and the actual optimizers for our problems appear to either have

only finitely many zeros in some cases (but not their Fourier transform). Indeed, imposing more zeros for f usually makes our bounds worse and the zeros have a tendency to fly off to infinity or collide when we run the optimizer.

The strategy we have described above is the one we use in low genus (or for small systole). In high genus (or for large systole), our asymptotic bounds are obtained by using special test functions related to Bessel functions and optimizing over certain parameters.

4.1. Certifying inequalities on intervals. Despite the numerical optimization used to produce our bounds, the end results are rigorous. The reason is that we work with rational zeros and polynomials with rational coefficients, so the linear systems involved are solved exactly over the rational numbers. The polynomials we get thus have actual double zeros rather than approximate ones.

To ascertain that $f(x) = p(x^2)e^{-x^2/2} \ge 0$ for every x in a given interval [a, b], we apply Sturm's theorem to count the number of distinct roots of p in that interval and make sure that there are no more than the number of imposed double zeros. This implies that neither p nor f changes sign on the interval, and it then suffices to check that p is strictly positive at some point or that its second derivative is strictly positive at a double zero.

We will also sometimes need to certify inequalities involving transcendental functions over intervals. In these cases, we approximate the transcendental functions with truncated Taylor series and apply Sturm's theorem to these approximation. The functions we consider either have positive or alternating Taylor coefficients, allowing us to know if the approximations are from below or above.

In some cases, we need to find the minimum of a function $h(x) = r(x)e^{-x/2}$ on an interval [a, b], where r is a polynomial. Since $h'(x) = (r'(x) - r(x)/2)e^{-x/2}$ we can use Sturm's theorem to verify that h has at most one critical point on the interval. If it has one, then it suffices to verify that r'(a) - r(a)/2 > 0 and r'(b) - r(b)/2 < 0 so that the critical point is a local maximum. The minimum of h is then at one of the endpoints, and this is also true if there is no critical point in the interval.

4.2. Certifying error bounds on integrals. Another difference with sphere packing bounds is that we have inequalities involving integrals that need to be checked. For example, one of our bounds requires that

(4.1)
$$\widehat{f}(i/2) \ge 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) \, dx$$

To obtain functions that satisfy this inequality, we first compute numerical approximations I_n of the integrals

$$\int_0^\infty L_n^{(-1/2)}(x^2) e^{-x^2/2} x \tanh(\pi x) \, dx$$

In our linear system of equations, we then impose that $\widehat{f}(i/2) = q(-1/4)e^{1/8}$ is equal to ρ times the numerical approximation of $2(g-1)\int_0^\infty \widehat{f}(x)x \tanh(\pi x) dx$ (given by a linear combination of the approximations I_n), where $\rho > 1$ is some rational number. Technically, we also replace $e^{1/8}$ by a rational approximation in this equation.

Once we have found a good candidate function f, we verify a posteriori that inequality (4.1) is satisfied. This is done by evaluating the left-hand side using interval arithmetic (which provides true lower and upper bounds on $\hat{f}(i/2)$) and finding certified bounds on the integral.

For a function h that is analytic in a neighborhood of a compact interval [a, b], the Arb package [Joh17] in SageMath [The21] is able to compute the integral $\int_a^b h(x) dx$ with certified error bounds. However, improper integrals (and in particular infinite intervals) cannot be handled. We thus use the Arb package to estimate $\int_0^b \hat{f}(x)x \tanh(\pi x) dx$ for some large b and then estimate the remainder $\int_b^\infty \hat{f}(x)x \tanh(\pi x) dx$ separately. For this, we use the inequalities

$$\tanh(\pi b) \le \tanh(\pi x) \le 1$$

for $x \ge b$. In all cases, our hypotheses will require that \hat{f} is eventually of constant sign, so it remains to estimate

$$\int_b^\infty \widehat{f}(x)x\,dx = \int_b^\infty xq(x^2)e^{-x^2/2}\,dx.$$

However, since $xq(x^2)$ is an odd polynomial, the function $xq(x^2)e^{-x^2/2}$ admits an explicit primitive and the integral can be computed exactly.

We will sometimes have to deal with more complicated integrals, in which case we estimate the remainder terms using ad hoc inequalities.

4.3. Ancillary files. Whenever we require certified error bounds on integrals in a proof, we explain how to estimate these integrals in the proof and state the resulting estimate that was obtained using interval arithmetic in SageMath. The calculations behind these estimates are all contained in the Jupyter notebook certified_integrals.ipynb attached as an ancillary file to the arXiv version of this paper.

Then there is one file verify_invariant.ipynb for each of the invariants we consider. Each such file contains a function invariant_poly which computes a pair of polynomials (p,q) such that $f(x) = p(x^2)e^{-x^2/2}$ and $\hat{f}(x) = q(x^2)e^{-x^2/2}$ are the Fourier transform of one another given a list of double zeros for each and perhaps additional data. Another function invariant_verify checks that all the required conditions on f and \hat{f} are satisfied and outputs a resulting rigorous upper bound on the invariant in question. The lists of input parameters that we used to produce the bounds in Tables 1 to 7 are stored in various files parameters_invariant.sobj that are loaded in the last cell of the verify_invariant notebook. Upon execution of this last cell, the program runs the invariant_verify function on each of these input parameters and prints out the resulting bounds.

5. Bessel functions

Bessel functions were used in [CE03] to obtain a new proof of the second best asymptotic upper bound on the density of sphere packings in \mathbb{R}^n due to Levenshtein [Lev79]. We will also use these functions to obtain our asymptotic bounds. We list some of their properties here for later reference.

One of the many equivalent definitions [Wat95, p.40] of the Bessel function of the first kind of order α is

$$J_{\alpha}(z) := \left(\frac{z}{2}\right)^{\alpha} \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \, \Gamma(n+\alpha+1)} \left(\frac{z}{2}\right)^{2n}$$

when α is not a negative integer, where Γ is the classical gamma function. For non-integer orders J_{α} is a multi-valued function, but by abuse of notation the quotient

$$\frac{J_{\alpha}(z)}{z^{\alpha}} = \frac{1}{2^{\alpha}} \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \, \Gamma(n+\alpha+1)} z^{2n}$$

defines an even entire function that takes the value $2^{-\alpha}/\Gamma(\alpha+1)$ at the origin. This leads us to define the *normalized Bessel function*

$$\eta_{\alpha}(z) := \Gamma(\alpha+1) \sum_{n=0}^{\infty} \frac{(-1)^n}{n! \,\Gamma(n+\alpha+1)} z^{2n} = 2^{\alpha} \Gamma(\alpha+1) \frac{J_{\alpha}(z)}{z^{\alpha}}$$

satisfying $\eta_{\alpha}(0) = 1$. For $\alpha > -1/2$, Poisson's integral formula for Bessel functions [Wat95, p.165] can be written as

$$\eta_{\alpha}(x) = \frac{1}{B(\frac{1}{2}, \alpha + \frac{1}{2})} \int_{0}^{1} (1 - t^{2})^{\alpha - 1/2} \cos(xt) dt$$

where

$$B(a,b) = \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)}$$

is the Beta function. This means that η_{α} is the Fourier transform of

$$\chi_{\alpha}(t) = \frac{\sqrt{\pi/2}}{B(\frac{1}{2}, \alpha + \frac{1}{2})} \operatorname{rect}(t/2)(1 - t^2)^{\alpha - 1/2}$$

where rect is the characteristic function of the interval $\left[-\frac{1}{2}, \frac{1}{2}\right]$. By Fourier inversion, we have

$$\widehat{\eta_{\alpha}}(t) = \chi_{\alpha}(t)$$

whenever $\alpha > 1/2$, which is when η_{α} is integrable. In particular, η_{α} is positive-definite for $\alpha > 1/2$ and its Fourier transform is supported in [-1, 1]. By the easy direction of the Paley–Wiener theorem, this implies that η_{α} has exponential type 1. In fact, along the imaginary axis we have the following exact asymptotic for every $\alpha \ge -1/2$ [Wat95, p.203]:

(5.1)
$$J_{\alpha}(ix) \sim \frac{e^x}{\sqrt{2\pi x}} \quad \text{as } x \to \infty \text{ in } \mathbb{R}.$$

The above integrability condition on ψ_{α} follows from the asymptotic formula

(5.2)
$$J_{\alpha}(z) = \sqrt{\frac{2}{\pi z}} \left(\cos\left(z - \frac{\alpha \pi}{2} - \frac{\pi}{4}\right) + O(e^{|\operatorname{Im} z|}/|z|) \right)$$

as $|z| \to \infty$ with $|\arg z| < \pi$ [AS64, p.364]. Also note that $J_{\alpha}(x)$ vanishes to order α at the origin, so that for $\alpha \ge -1/2$ the function $x \mapsto \sqrt{x} J_{\alpha}(x)$ is bounded near the origin and hence on $(0, \infty)$ by continuity and the above asymptotic.

We will frequently make use of the even entire functions

$$\varphi_{\alpha}(z) = \frac{J_{\alpha}(z/2)^2}{z^{2\alpha}} = \left(\frac{\eta_{\alpha}(z/2)}{4^{\alpha}\Gamma(\alpha+1)}\right)^2$$

and

$$\psi_{\alpha}(z) = \frac{J_{\alpha}(z)^2}{z^{2\alpha}(1 - (z/j_{\alpha})^2)}$$

where j_{α} is the first positive root of J_{α} . These are such that $\varphi_{\alpha}(x) \geq 0$ for every $x \in \mathbb{R}$ and $\psi_{\alpha}(x) \leq 0$ for $|x| \geq j_{\alpha}$. Up to positive constants, $\varphi_{\alpha}(2x)$ is equal to $\widehat{\chi_{\alpha} * \chi_{\alpha}}(x)$ so its Fourier transform is $\chi_{\alpha} * \chi_{\alpha} \geq 0$ as long as φ_{α} and χ_{α} are integrable, which holds whenever $\alpha > 0$. In other words, φ_{α} is positive-definite if $\alpha > 0$, with Fourier transform supported in [-1, 1]. It was also shown in [GIT20, Remark 1.1] that ψ_{α} is positive-definite if $\alpha \geq -1/2$. Observe that $\widehat{\varphi_{\alpha}}$ is admissible if $\alpha > 1/2$ and $\widehat{\psi_{\alpha}}$ is admissible if $\alpha > -1/2$ by the asymptotic formula (5.2).

6. Systole

6.1. The criterion. The systole of a closed hyperbolic surface is defined as the length of any of its shortest closed geodesics (also called systoles). Our criterion for bounding the systole goes as follows.

Theorem 6.1. Let $g \ge 2$. Suppose that f is a non-constant admissible function for which there exists an R > 0 such that

• $f(x) \le 0$ if $x \ge R$;

•
$$\widehat{f}(\xi) \ge 0$$
 for every $\xi \in \mathbb{R} \cup i\left[-\frac{1}{2}, \frac{1}{2}\right]$;

• $\widehat{f}(i/2) \ge 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) dx.$

Then $sys(M) \leq R$ for every closed hyperbolic surface M of genus g.

Proof. Suppose that there is a hyperbolic surface M of genus g such that $\operatorname{sys}(M) > R$. Then $\operatorname{sys}(N) > R$ for every surface N in some connected neighborhood U of M in moduli space. This implies that $f(\ell(\gamma)) \leq 0$ for every $\gamma \in \mathcal{C}(N)$ and we also have $\widehat{f}\left(\sqrt{\lambda_j(N) - \frac{1}{4}}\right) \geq 0$ for every $j \geq 0$ by the hypotheses on f and \widehat{f} . From the Selberg trace formula, we obtain

$$\begin{aligned} \widehat{f}(i/2) &\leq \sum_{j=0}^{\infty} \widehat{f}\left(\sqrt{\lambda_j(N) - \frac{1}{4}}\right) \\ &= 2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(N)} \frac{\Lambda(\gamma)f(\ell(\gamma))}{2\sinh(\ell(\gamma)/2)} \\ &\leq 2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx \\ &\leq \widehat{f}(i/2) \end{aligned}$$

for every $N \in U$. We conclude that $\widehat{f}\left(\sqrt{\lambda_j(N) - \frac{1}{4}}\right) = 0$ for every $j \ge 1$. Since \widehat{f} is holomorphic in a strip and not constant equal to zero, its zeros are isolated. This implies that for every $j \ge 1$, the eigenvalue $\lambda_j(N)$ is a constant function of $N \in U$ since eigenvalues depend continuously on the metric (see e.g. [BU83]). Therefore, all the surfaces in U are isospectral. However, Gel'fand proved that any continuous deformation of M that preserves the entire Laplace spectrum is constant [Gel63], which is a contradiction.

Remark 6.2. The analogous result for flat tori was proved in [CE03, Theorem 3.2] using a rescaling and limiting argument for the second half of the proof.

Remark 6.3. It is easy to see that if the inequality in the third bullet point is strict, then the conclusion can be strengthened to a strict inequality. The proof proceeds similarly as above, but the chain of inequalities directly leads to a contradiction.

6.2. Low genus. The upper bounds we have obtained from Theorem 6.1 though numerical optimization are listed in Table 1 for $2 \le g \le 20$. The verification of these values is done in the ancillary file verify_systole.ipynb. They are lower than the previous best upper bounds except in genus 2 where the optimal bound is $2 \operatorname{arccosh}(1+\sqrt{2}) \approx 3.057142$ [Jen84]. In all other genera, the previous best upper bound was Bavard's inequality [Bav96]

(6.1)
$$\operatorname{sys}(M) \le 2 \operatorname{arccosh}\left(\frac{1}{2\sin(\pi/(12g-6))}\right),$$

which comes from a sharp upper bound on the radius of an embedded disk in M.

We have also listed the largest recorded value of the systole in some genera. Those listed in genus 7, 14, and 17 come from Hurwitz surfaces. Technically, the values from [Vog03] and [SS22] were obtained by numerical calculations in triangle groups and are not completely rigorous, but they could be made rigorous in principle (this was done in [DT00] for the Klein quartic and in [Woo01] for the Hurwitz surfaces of genus 14). For Hurwitz surfaces, the calculations from [SS22] corroborate those of [Vog03].

Since the systole does not decrease under covers, one could fill in all the blanks in the table with values in lower genera. Similarly, the value listed in genus 13 persists in every genus g > 13 [FBR22]. We decided not to list these since better constructions surely exist.

genus	lower bound	LP bound	previous upper bound
2	3.057141 [Jen84]	3.156053^{\dagger}	3.057142 [Jen84]
3	3.983304 [Sch93]	4.194719	4.494373 [Bav96]
4	4.624499 [Sch93]	4.876863	5.176481 [Bav96]
5	4.91456 [Sch94]	5.381937	5.682841 [Bav96]
6	5.109 [CB05]	5.783671	6.086062 [Bav96]
7	5.796298 [Vog03]	6.117160	6.421249 [Bav96]
8		6.407734	6.708126 [Bav96]
9	5.376 [Sch93]	6.655635	6.958903 [Bav96]
10		6.880869	7.181671 [Bav96]
11	5.980406 [Sch93]	7.080715	7.382068 [Bav96]
12		7.262735	7.564184 [Bav96]
13	5.909039 [FBR22]	7.429527	7.731080 [Bav96]
14	6.887905 [Woo01, Vog03]	7.584859	7.885106 [Bav96]
15		7.729299	8.028108 [Bav96]
16		7.863529	8.161558 [Bav96]
17	7.609407 [Vog03]	7.988773	8.286655 [Bav96]
18		8.118854	8.404383 [Bav96]
19	7.358 [SS22]	8.220710	8.515562 [Bav96]
20		8.328393	8.620882 [Bav96]

TABLE 1. Bounds on the maximum of sys. The dagger indicates where the linear programming bound fails to beat the previous best upper bound.

6.3. Asymptotics. Note that the term $\sin(\pi/(12g-6))$ appearing in Bavard's bound is asymptotic to $\frac{\pi}{12g}$ as $g \to \infty$ and

$$\operatorname{arccosh}(x) = \log(x + \sqrt{x^2 - 1}) = \log(x) + \log(2) + o(1)$$

as $x \to \infty$ so that Bavard's bound (6.1) can be rewritten as

$$sys(M) \le 2\log\left(\frac{6g}{\pi}\right) + 2\log(2) + o(1) = 2\log(g) + 2\log\left(\frac{12}{\pi}\right) + o(1) = 2\log(g) + 2.680353... + o(1)$$

as $g \to \infty$. By comparison, the elementary area bound coming from the fact that a disk of radius sys(M)/2 is embedded is

$$sys(M) \le 4 \operatorname{arcsinh}(\sqrt{g-1}) \\ = 2\log(g) + 4\log(2) + o(1) \\ = 2\log(g) + 2.772588 \dots + o(1)$$

as $g \to \infty$.

We will decrease the additive constant in Bavard's bound by roughly 0.271, which is consistent with the improvement we have observed in small genus.

Theorem 6.4. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$sys(M) < 2\log(g) + 2.409.$$

Remark 6.5. In terms of area, this result means that in large genus, a disk of radius sys(M)/2 cannot occupy a proportion of more than

$$\frac{e^{2.409/2}}{4} \approx 0.833772\dots$$

of M, while a maximal embedded disk can occupy as much as

$$\frac{3}{\pi} \approx 0.954929\dots$$

of the surface by Bavard's result [Bav96].

Remark 6.6. In large genus, the constructions with the fastest growing systole known are given by towers of principal congruence covers of arithmetic surfaces. For each such tower, there is a constant c such that

$$\operatorname{sys}(M) \ge \frac{4}{3}\log(g) - c$$

for every surface M of genus g in the tower [KSV07, Theorem 1.5].

The lengthy proof of Theorem 6.4 will require several estimates presented in the form of lemmata below. The strategy is to apply Theorem 6.1 with functions f such that

$$\widehat{f}(x) = h_c(bx)\varphi_\alpha(Rx),$$

for some parameters α , b, c and R, where

$$h_c(x) = (c - 1 + x^2)e^{-x^2/2}$$
 and $\varphi_{\alpha}(x) = \frac{J_{\alpha}(x/2)^2}{x^{2\alpha}}$

with J_{α} the Bessel function of order α as in Section 5. We were led to these types of functions by studying the numerical data gathered in small genus. We believe they are

nearly optimal. Using functions of the form $\hat{f}(x) = \varphi_{\alpha}(Rx)$ instead yields a bound of the form $2\log(g) + c$ but with a c larger that Bavard's, which is why we need to use more complicated functions.

The parameters α, b, c will be fixed at some point and only R will depend on the genus. Since we need

$$\widehat{f}(i/2) > 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) \, dx \ge 0.$$

we require that

$$c - 1 - b^2/4 > 0$$
 or equivalently $c > 1 + b^2/4$,

which in turn implies that \widehat{f} is non-negative on $\mathbb{R} \cup i[-\frac{1}{2}, \frac{1}{2}]$.

To apply Theorem 6.1, we need to show that f is eventually negative and to estimate the location of its last sign change. By Fourier inversion and the convolution theorem, we have

(6.2)
$$f = \frac{1}{\sqrt{2\pi}} \, \widehat{h_c^b} * \widehat{\varphi_\alpha^R}$$

where $h_c^b(x) = h_c(bx)$ and $\varphi_{\alpha}^R(x) = \varphi_{\alpha}(Rx)$. Recall that φ_{α} is positive-definite with Fourier transform supported in [-1, 1] provided that $\alpha > 0$, as explained in Section 5. It follows that $\widehat{\varphi_{\alpha}^R}(x) = \widehat{\varphi_{\alpha}}(x/R)/R$ is non-negative and supported in [-R, R] while

$$\widehat{h_c^b}(x) = \frac{1}{b}\widehat{h_c}(x/b) = \frac{1}{b}(c - x^2/b^2)e^{-x^2/(2b^2)}$$

is non-positive outside $[-b\sqrt{c}, b\sqrt{c}]$. From this, it is easy to show that the convolution f is non-positive outside $[-R - b\sqrt{c}, R + b\sqrt{c}]$. However, $f(R + b\sqrt{c}) < 0$ so that the last sign change occurs before that. Here is how we can estimate its location more precisely.

Lemma 6.7. Suppose that $\kappa_0 \in \mathbb{R}$ is such that

$$\int_{\kappa_0}^{\infty} (x - \kappa_0)^{2\alpha} (c - x^2) e^{-x^2/2} \, dx < 0$$

and let f be defined as above with $\alpha \in (0,1)$, b > 0, and c > 0. If $\kappa_n \to \kappa_0$ and $R_n \to \infty$ as $n \to \infty$, then $R_n^{2\alpha+1}f(R_n + b\kappa_n)$ converges to some limit $\mu_{\alpha,b,c,\kappa_0} < 0$ as $n \to \infty$, where $\mu_{\alpha,b,c,\kappa_0}$ depends continuously on the parameters. In particular, $f(R_n + b\kappa_n)$ is negative whenever n is large enough.

The proof will require the following lemma.

Lemma 6.8. For every $\alpha \in (0,1)$, there exist constants $c_{\alpha}, d_{\alpha} > 0$ such that

$$R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-\frac{y}{R}\right) \le d_{\alpha}y^{2\alpha}$$

for every R > 0 and y > 0 and

$$\lim_{R \to \infty} R^{2\alpha} \widehat{\varphi_{\alpha}} \left(1 - \frac{y}{R} \right) = c_{\alpha} y^{2\alpha}$$

for every y > 0. Moreover, c_{α} depends continuously on α .

Taking this lemma for granted for a moment, let us prove the preceding one.

Proof of Lemma 6.7. We have

$$\begin{split} R^{2\alpha+1}f(R+b\kappa) &= \frac{1}{\sqrt{2\pi}}R^{2\alpha+1}\int_{-\infty}^{\infty}\widehat{h_c^b}(x)\widehat{\varphi_{\alpha}^R}(R+b\kappa-x)\,dx\\ &= \frac{1}{b\sqrt{2\pi}}R^{2\alpha}\int_{-\infty}^{\infty}\widehat{h_c}(x/b)\widehat{\varphi_{\alpha}}\left(1-\frac{x-b\kappa}{R}\right)dx\\ &= \frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty}\widehat{h_c}(y)R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-b\frac{y-\kappa}{R}\right)dy\\ &= \frac{1}{\sqrt{2\pi}}\int_{\kappa}^{\infty}\widehat{h_c}(y)R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-b\frac{y-\kappa}{R}\right)dy \end{split}$$

by the change of variable y = x/b, where the last equality is because $\widehat{\varphi_{\alpha}}$ vanishes after 1. If $y > \kappa$, then $R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1 - b\frac{y-\kappa}{R}\right) \leq d_{\alpha}(b(y-\kappa))^{2\alpha}$ for some $d_{\alpha} > 0$ according to

If $y > \kappa$, then $R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-b\frac{y-\kappa}{R}\right) \leq d_{\alpha}(b(y-\kappa))^{2\alpha}$ for some $d_{\alpha} > 0$ according to Lemma 6.8. If we write $\underline{\kappa} = \inf_{n} \kappa_{n}$, then we have that $\widehat{h_{c}}(y)R_{n}^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-b\frac{y-\kappa_{n}}{R_{n}}\right)$ is bounded independently of n by the integrable function

$$d_{\alpha}b^{2\alpha}(y-\underline{\kappa})^{2\alpha}|\widehat{h_c}(y)|$$

on the interval $[\underline{\kappa}, \infty)$. We can therefore apply the dominated convergence theorem to conclude that

$$\lim_{R_n \to \infty} R_n^{2\alpha+1} f(R_n + b\kappa_n) = \frac{c_\alpha b^{2\alpha}}{\sqrt{2\pi}} \int_{\kappa_0}^\infty (y - \kappa_0)^{2\alpha} \widehat{h_c}(y) \, dy < 0,$$

where we used the limit from Lemma 6.8 and the hypothesis that the last integral is negative. It follows that $f(R_n + b\kappa_n)$ is eventually negative. That the limit depends continuously on the parameters is a consequence of the dominated convergence theorem.

We now prove the lemma about the behaviour of $\widehat{\varphi_{\alpha}}$ near the end of its support.

Proof of Lemma 6.8. Since $\widehat{\varphi_{\alpha}}$ is continuous (because φ_{α} is integrable if $\alpha > 0$) and supported in [-1, 1], we have $\widehat{\varphi_{\alpha}}(1) = 0$ and

$$\begin{aligned} \widehat{\varphi_{\alpha}} \left(1 - \frac{y}{R} \right) &= \widehat{\varphi_{\alpha}} \left(1 - \frac{y}{R} \right) + \widehat{\varphi_{\alpha}} \left(1 + \frac{y}{R} \right) - 2\widehat{\varphi_{\alpha}}(1) \\ &= \sqrt{\frac{2}{\pi}} \int_{0}^{\infty} \varphi_{\alpha}(x) \left(\cos \left(x - \frac{xy}{R} \right) + \cos \left(x + \frac{xy}{R} \right) - 2\cos(x) \right) dx. \\ &= 2\sqrt{\frac{2}{\pi}} \int_{0}^{\infty} \varphi_{\alpha}(x) \cos(x) \left(\cos \left(\frac{xy}{R} \right) - 1 \right) dx \\ &= R\sqrt{\frac{8}{\pi}} \int_{0}^{\infty} \varphi_{\alpha}(Ru) \cos(Ru) \left(\cos(yu) - 1 \right) du \end{aligned}$$

by the change of variable u = x/R.

Observe that

$$R^{2\alpha+1}\varphi_{\alpha}(Ru) = \frac{1}{u^{2\alpha+1}}RuJ_{\alpha}(Ru/2)^2 \le \frac{D_{\alpha}}{u^{2\alpha+1}}$$

for some constant $D_{\alpha} > 0$ since the function $x J_{\alpha}(x/2)^2$ is bounded on the positive real axis (see Section 5). It follows that

$$R^{2\alpha}\widehat{\varphi_{\alpha}}\left(1-\frac{y}{R}\right) \le D_{\alpha}\sqrt{\frac{8}{\pi}} \int_{0}^{\infty} \frac{1-\cos(yu)}{u^{2\alpha+1}} \, du = y^{2\alpha}D_{\alpha}\sqrt{\frac{8}{\pi}} \int_{0}^{\infty} \frac{1-\cos(v)}{v^{2\alpha+1}} \, dv$$

where v = yu. The first assertion is thus proved with

$$d_{\alpha} = D_{\alpha} \sqrt{\frac{8}{\pi}} \int_0^\infty \frac{1 - \cos(v)}{v^{2\alpha + 1}} \, dv,$$

which is finite because of the hypothesis on α . Indeed, $2\alpha - 1 < 1$ implies that

$$G_{\alpha}(v) = \frac{1 - \cos(v)}{v^{2\alpha + 1}}$$

is integrable near the origin and $2\alpha + 1 > 1$ implies that it is integrable near infinity.

For the second assertion, we will use the asymptotic expansion of Bessel functions along the real line. We have

$$(Ru)^{2\alpha+1}\varphi_{\alpha}(Ru) = RuJ_{\alpha}(Ru/2)^{2}$$

$$= \frac{4}{\pi} \left(\cos\left(\frac{Ru}{2} - \frac{(2\alpha+1)\pi}{4}\right) + O\left(\frac{1}{Ru}\right) \right)^{2}$$

$$= \frac{4}{\pi} \left(\cos^{2}\left(\frac{Ru}{2} - \frac{(2\alpha+1)\pi}{4}\right) + O\left(\frac{1}{Ru}\right) \right).$$

Moreover, recall that $(Ru)^{2\alpha+1}\varphi_{\alpha}(Ru)$ is uniformly bounded so that

$$R^{2\alpha+1}\varphi_{\alpha}(Ru)(1-\cos(yu)) = (Ru)^{2\alpha+1}\varphi_{\alpha}(Ru)G_{\alpha}(yu)y^{2\alpha+1}$$

is bounded by an integrable function and similarly for

$$F_R(u) = \left(R^{2\alpha+1} \varphi_\alpha(Ru) - \frac{4}{\pi u^{2\alpha+1}} \cos^2\left(\frac{Ru}{2} - \frac{(2\alpha+1)\pi}{4}\right) \right) \cos(Ru)(1 - \cos(yu)),$$

namely, by some constant multiple of $G_{\alpha}(yu)$. Furthermore, for every u > 0 we have that $F_R(u)$ tends to zero as $R \to \infty$ by the asymptotic expansion above. By the dominated convergence theorem, $\int_0^{\infty} F_R(u) du \to 0$ as $R \to \infty$. We conclude that

$$\lim_{R \to \infty} \int_0^\infty R^{2\alpha + 1} \varphi_\alpha(Ru) \cos(Ru) \left(\cos(yu) - 1 \right) \, du$$

is equal to

$$-\frac{4y^{2\alpha+1}}{\pi}\lim_{R\to\infty}\int_0^\infty\cos^2\left(\frac{Ru}{2}-\frac{(2\alpha+1)\pi}{4}\right)\cos(Ru)G_\alpha(yu)\,du$$

provided that either limit exists. Using the identity

$$\cos^{2}(a-b) = \frac{1+\cos(2(a-b))}{2} = \frac{1}{2} \left(1+\cos(2a)\cos(2b)+\sin(2a)\sin(2b)\right),$$

we find that the integral inside the second limit is equal to

$$\frac{1}{2}\int_0^\infty \left(1 + \cos(Ru)\cos\left(\frac{(2\alpha+1)\pi}{2}\right) + \sin(Ru)\sin\left(\frac{(2\alpha+1)\pi}{2}\right)\right)\cos(Ru)G_\alpha(yu)du,$$

which we split into a sum of three terms. The first and last terms tend to zero as $R \to \infty$ by the Riemann–Lebesgue lemma, where we used the identity $2\cos(Ru)\sin(Ru) = \sin(2Ru)$. Writing $\cos^2(Ru) = \frac{1+\cos(2Ru)}{2}$ and applying the Riemann–Lebesgue lemma again shows that the middle term converges to

$$\frac{1}{4}\cos\left(\frac{(2\alpha+1)\pi}{2}\right)\int_0^\infty G_\alpha(yu)\,du = \frac{1}{4y}\cos\left(\frac{(2\alpha+1)\pi}{2}\right)\int_0^\infty G_\alpha(v)\,dv.$$

The result then follows with

$$c_{\alpha} = -\left(\frac{2}{\pi}\right)^{3/2} \cos\left(\frac{(2\alpha+1)\pi}{2}\right) \int_0^\infty G_{\alpha}(v) \, dv$$

which is a positive number since $\frac{(2\alpha+1)\pi}{2} \in (\pi, \frac{3\pi}{2})$ and G_{α} is integrable and non-negative. That c_{α} depends continuously on α follows from the continuity of the cosine function, the continuous dependence of G_{α} on α , and the dominated convergence theorem.

We now need to check that the hypothesis of Lemma 6.7 is satisfied for some specific parameters.

Lemma 6.9. Let $\alpha = 0.559$ and c = 2.3726. Then

$$\int_{\kappa}^{\infty} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx < 0.$$

for every $\kappa \geq \kappa_0 = 0.1814$.

Proof. If $\kappa \geq \sqrt{c}$, then the result is obvious, because the integrand is non-positive. So we concentrate on the interval $[\kappa_0, \sqrt{c}]$.

Let us denote by the integral in the statement of the lemma by $I(\kappa)$. We first verify that $I(\kappa_0) < 0$ using interval arithmetic in SageMath. Since the integrand has a singularity at κ_0 , we need to estimate the integral differently near there. We write

$$\int_{\kappa_0}^{\kappa_0+A} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx \le A^{2\alpha} \int_{\kappa_0}^{\kappa_0+A} (c-x^2) e^{-x^2/2} \, dx$$

and then

$$\int_{\kappa_0+A}^{\infty} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx \le \int_{\kappa_0+A}^{B} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx$$

as long as A > 0, $B \ge \kappa_0 + A$, and $B \ge \sqrt{c}$. With $A = 10^{-4}$ and B = 10, these estimates provide the certified upper bound

 $I(\kappa_0) \le -0.0000117812526025449 < 0.$

We then show that $I'(\kappa) < 0$ on $[\kappa_0, 0.7]$ as follows. By the change of variable $y = x - \kappa$ we get

$$I(\kappa) = \int_0^\infty y^{2\alpha} (c - (y + \kappa)^2) e^{-(y + \kappa)^2/2} \, dy$$

and then differentiation under the integral sign (which is justified because the derivative of the integrand is uniformly bounded by a polynomial times a Gaussian for κ in a bounded

interval) gives

$$\begin{split} I'(\kappa) &= \int_0^\infty y^{2\alpha} (y+\kappa) ((y+\kappa)^2 - (c+2)) e^{-(y+\kappa)^2/2} \, dy \\ &= \int_\kappa^\infty (x-\kappa)^{2\alpha} x (x^2 - (c+2)) e^{-x^2/2} \, dx \\ &= \int_\kappa^{\sqrt{c+2}} (x-\kappa)^{2\alpha} x (x^2 - (c+2)) e^{-x^2/2} \, dx + \int_{\sqrt{c+2}}^\infty (x-\kappa)^{2\alpha} x (x^2 - (c+2)) e^{-x^2/2} \, dx \\ &\leq \int_{0.7}^{\sqrt{c+2}} (x-0.7)^{2\alpha} x (x^2 - (c+2)) e^{-x^2/2} \, dx + \int_{\sqrt{c+2}}^\infty (x-\kappa_0)^{2\alpha} x (x^2 - (c+2)) e^{-x^2/2} \, dx \end{split}$$

whenever $\kappa_0 \leq \kappa \leq 0.7$. We verify that this sum of integrals is at most

-0.0464961225743898

(hence negative) using interval arithmetic (again splitting the integrals near 0.7 and ∞). It follows that I is bounded above by $I(\kappa_0) < 0$ on $[\kappa_0, 0.7]$.

For any $\kappa \in [0.7, \sqrt{c}]$, we estimate

$$I(\kappa) = \int_{\kappa}^{\sqrt{c}} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} dx + \int_{\sqrt{c}}^{\infty} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} dx$$
$$\leq \int_{0.7}^{\sqrt{c}} (x-0.7)^{2\alpha} (c-x^2) e^{-x^2/2} dx + \int_{\sqrt{c}}^{\infty} (x-\sqrt{c})^{2\alpha} (c-x^2) e^{-x^2/2} dx$$
$$\leq -0.0907427113682867$$

using interval arithmetic once again, which completes the proof.

From the above lemma, we can deduce that the function f from equation (6.2) is nonpositive from $R + b\kappa_0$ onwards provided that R is large enough.

Corollary 6.10. Let $\alpha = 0.559$, b > 0, c = 2.3726, and $\kappa_0 = 0.1814$. Then there exists some $R_0 > 0$ such that $f(R + b\kappa) \leq 0$ for every $R \geq R_0$ and every $\kappa \geq \kappa_0$, where f is as in equation (6.2).

Proof. If R > 0 and $\kappa \ge \sqrt{c}$, then

$$f(R+b\kappa) = \frac{1}{\sqrt{2\pi}} \int_{\kappa}^{\infty} \widehat{h_c}(y)\widehat{\varphi_{\alpha}}\left(1 - b\frac{y-\kappa}{R}\right) dy \le 0$$

because $\widehat{h_c}$ is non-positive after \sqrt{c} and $\widehat{\varphi_{\alpha}}$ is non-negative, so their product is non-positive.

For every $\kappa \in [\kappa_0, \sqrt{c}]$ we have that $R^{2\alpha+1}f(R+b\kappa)$ converges to some $\mu_{\alpha,b,c,\kappa} < 0$ as $R \to \infty$ by Lemma 6.7 and Lemma 6.9. Let

$$F(R,\kappa) := \begin{cases} R^{2\alpha+1}f(R+b\kappa) & \text{if } R \in (0,\infty) \\ \mu_{\alpha,b,c,\kappa} & \text{if } R = \infty \end{cases}$$

It follows form Lemma 6.7 that F is continuous at every point in $\{\infty\} \times [\kappa_0, \sqrt{c}]$ and the continuity on $(0, \infty) \times [\kappa_0, \sqrt{c}]$ is a consequence of the dominated convergence theorem.

We deduce that every $\kappa \in [\kappa_0, \sqrt{c}]$, there exists some neighborhood U of κ and some $R_{\kappa} > 0$ such that $R^{2\alpha+1}f(R+bu) < 0$ for every $u \in U$ and every $R \ge R_{\kappa}$. By compactness, we can find an $R_0 > 0$ that works for the whole interval $[\kappa_0, \sqrt{c}]$.

Armed with this result, we can finally give the proof of Theorem 6.4.

Proof of Theorem 6.4. Let $\hat{f}(x) = h_c(bx)\varphi_\alpha(Rx)$ as before with $\alpha = 0.559$, b = 1.0286, c = 2.3726, and R > 0. Note that $c > 1 + b^2/4$ as required, which implies that $\hat{f} \ge 0$ on $\mathbb{R} \cup i \begin{bmatrix} \frac{1}{2}, \frac{1}{2} \end{bmatrix}$. By Corollary 6.10, there exists some $R_0 > 0$ such that if $R \ge R_0$, then $f(x) \le 0$ whenever $x \ge R + b\kappa_0$, where $\kappa_0 = 0.1814$ (recall that f itself depends on R, which is why we need to consider the parameters $\kappa \ge \kappa_0$ to cover every $x \ge R + b\kappa_0$). Therefore, $R + b\kappa_0$ provides a bound on sys(M) as long as

$$\widehat{f}(i/2) \ge 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) \, dx.$$

We will thus estimate both sides and choose R so that this inequality is satisfied.

For the left-hand side, we have

$$\widehat{f}(i/2) = h_c(bi/2)\varphi_\alpha(iR/2)$$

where

$$\varphi_{\alpha}(iR/2) = \frac{J_{\alpha}(iR/4)^2}{(R/2)^{2\alpha}} \sim \frac{2^{2\alpha+1}}{\pi} \frac{e^{R/2}}{R^{2\alpha+1}}$$

as $R \to \infty$ by the asymptotic (5.1).

As for the integral term, we consider

$$\int_0^\infty h_c(bx)x\tanh(\pi x)R^{2\alpha+1}\varphi_\alpha(Rx)\,dx$$

and recall that for every x > 0 we have

$$\left| R^{2\alpha+1}\varphi_{\alpha}(Rx) - \frac{4}{\pi x^{2\alpha+1}}\cos^2\left(\frac{Rx}{2} - \frac{(2\alpha+1)\pi}{4}\right) \right| \to 0$$

as $R \to \infty$ and by a similar argument as in the proof of Lemma 6.8, we have

$$\lim_{R \to \infty} R^{2\alpha+1} \int_0^\infty h_c(bx) x \tanh(\pi x) \varphi_\alpha(Rx) \, dx$$
$$= \frac{4}{\pi} \lim_{R \to \infty} \int_0^\infty \frac{h_c(bx) \tanh(\pi x)}{x^{2\alpha}} \cos^2\left(\frac{Rx}{2} - \frac{(2\alpha+1)\pi}{4}\right) \, dx.$$

Here observe that $\frac{h_c(bx) \tanh(\pi x)}{x^{2\alpha}}$ is integrable as long as $2\alpha - 1 < 1$ or $\alpha < 1$, which is what is needed to apply the dominated convergence and obtain this equality. To compute the integral inside the limit, we again write

$$\cos^2\left(\frac{Rx}{2} - \frac{(2\alpha+1)\pi}{4}\right) = \frac{1}{2}\left(1 + \cos(Rx)\cos\left(\frac{(2\alpha+1)\pi}{2}\right) + \sin(Rx)\sin\left(\frac{(2\alpha+1)\pi}{2}\right)\right)$$

By the Riemann–Lebesgue lemma, the integrals of the terms with $\cos(Rx)$ or $\sin(Rx)$ tend to zero as $R \to \infty$ so that

$$\frac{4}{\pi} \lim_{R \to \infty} \int_0^\infty \frac{h_c(bx) \tanh(\pi x)}{x^{2\alpha}} \cos^2\left(\frac{Rx}{2} - \frac{(2\alpha + 1)\pi}{4}\right) dx = \frac{2}{\pi} \int_0^\infty \frac{h_c(bx) \tanh(\pi x)}{x^{2\alpha}} dx.$$

We thus have

$$\frac{f(i/2)}{\int_0^\infty \widehat{f}(x)x\tanh(\pi x)\,dx} \sim \frac{4^\alpha h_c(bi/2)}{\int_0^\infty \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}}dx} e^{R/2}$$

as $R \to \infty$.

For any $\rho > 1$ and $g \ge 2$, if we choose R such that the right-hand side equal is equal to $2(q-1)\rho$, then the left-hand side will be larger than 2(q-1) provided that R is large enough so that the asymptotic is sufficiently precise. We thus take

$$R = 2\log(g-1) + 2\log\left(\rho \frac{2^{1-2\alpha}}{h_c(bi/2)} \int_0^\infty \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}} dx\right),$$

which tends to infinity as $q \to \infty$. The hypotheses of Theorem 6.1 will then satisfied if q is large enough, and we can further assume that $R \geq R_0$, so that the resulting bound on the systole is

$$\operatorname{sys}(M) \le R + b\kappa_0$$

$$= 2\log(g-1) + 2\log(\rho) + b\kappa_0 + 2\log\left(\frac{2^{1-2\alpha}}{h_c(bi/2)}\int_0^\infty \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}}dx\right)$$

according to Corollary 6.10. Since $\rho > 1$ was arbitrary, the additive constant can be taken as close as we wish to

$$b\kappa_0 + 2\log\left(\frac{2^{1-2\alpha}}{h_c(bi/2)}\int_0^\infty \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}}dx\right)$$

To estimate the integral rigorously, we restrict to a compact interval $[A, B] \subset (0, \infty)$ and estimate the remaining parts by

$$\int_{(0,\infty)\setminus[A,B]} \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}} dx \le \pi \int_0^A (c-1+b^2A^2)x^{1-2\alpha} dx + \int_B^\infty \frac{h_c(bx)}{x^{2\alpha}} dx \le \pi (c-1+b^2A^2)\frac{A^{2-2\alpha}}{2-2\alpha} + \frac{h_c(bB)}{(2\alpha-1)B^{2\alpha-1}},$$

as long as B is large enough so that h_c is decreasing on $[bB, \infty)$. Since

$$h'_c(x) = x((3-c) - x^2)e^{-x^2/2}$$

is negative when $x > \sqrt{3-c}$ and since our parameters satisfy $\sqrt{3-c} < 1 < b$, any $B \ge 1$ works.

With $A = 10^{-6}$ and B = 10, interval arithmetic in SageMath certifies that

$$b\kappa_0 + 2\log\left(\frac{2^{1-2\alpha}}{h_c(bi/2)}\int_0^\infty \frac{h_c(bx)\tanh(\pi x)}{x^{2\alpha}}dx\right) \le 2.40896511079437 < 2.409$$
given parameters, as required.

at the given parameters, as required.

7. Kissing numbers

7.1. The criterion. The kissing number of a hyperbolic surface M is defined as its number of oriented closed geodesics of minimal length. In the literature it is common to consider the quantity $\frac{1}{2}$ kiss(M), which counts the number of unoriented systoles in M. Our criterion for bounding kissing numbers is the same as the one we used in [FBP22] for hyperbolic manifolds in any dimension. We repeat the proof in the surface case for convenience.

Theorem 7.1. Let s > 0 and suppose that f is an admissible function such that

•
$$\widehat{f}(\xi) \ge 0$$
 for every $\xi \in \mathbb{R} \cup i \left[-\frac{1}{2}, \frac{1}{2}\right]$;

• $f(x) \leq 0$ whenever $x \geq s$ and f(s) < 0.

Then for every closed hyperbolic surface M of genus $g \ge 2$ such that sys(M) = s we have

kiss
$$(M) \le 4\sqrt{2\pi}(g-1)\frac{\sinh(s/2)}{s|f(s)|} \int_0^\infty \widehat{f}(x)x\tanh(\pi x)\,dx.$$

Proof. We have

$$0 \leq \sum_{j=0}^{\infty} \widehat{f}\left(\sqrt{\lambda_j(M)} - \frac{1}{4}\right)$$

= $2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(M)} \frac{\Lambda(\gamma)}{2\sinh(\ell(\gamma)/2)} f(\ell(\gamma))$
 $\leq 2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx + \frac{\operatorname{kiss}(M)}{\sqrt{2\pi}} \frac{\operatorname{sys}(M)}{2\sinh(\operatorname{sys}(M)/2)} f(\operatorname{sys}(M))$

since the contribution of geodesics longer than the systole is non-positive. Rearranging gives the desired result. $\hfill \Box$

Remark 7.2. We could take the value of $\hat{f}(\sqrt{-1/4}) = \hat{f}(i/2)$ into account to get an a priori better bound, but in practice we have found that the optimal functions seem to satisfy $\hat{f}(i/2) = 0$. We therefore enforce this condition in our program to speed up convergence.

Definition 7.3. For s > 0, we define K(s) as the infimum of

$$4\sqrt{2\pi}\frac{\sinh(s/2)}{s|f(s)|}\int_0^\infty \widehat{f}(x)x\tanh(\pi x)\,dx$$

over the functions f that satisfy the hypotheses of Theorem 7.1.

Theorem 7.1 can then be restated as saying that

$$kiss(M) \le K(sys(M))(g-1).$$

The following monotonicity result will greatly simplify our task of finding a global bound (independent of the systole) in each genus.

Lemma 7.4. The function K(s) is non-decreasing for $s \in [6, \infty)$.

Proof. Let f be a function that satisfies the hypotheses of Theorem 7.1 at some $s_2 > 6$ and let $s_1 \in [6, s_2)$. We then consider the function $\phi(x) = f\left(\frac{s_2}{s_1}x\right)$ with $\hat{\phi}(x) = \frac{s_1}{s_2}\hat{f}\left(\frac{s_1}{s_2}x\right)$. By the hypotheses on f we have that $\phi(x) = f\left(\frac{s_2}{s_1}x\right) \leq 0$ whenever $x \geq s_1$, because that implies $\frac{s_2}{s_1}x \geq s_2$. Moreover, we have $\phi(s_1) = f(s_2) < 0$. Secondly, if $\xi \in \mathbb{R} \cup i \left[-\frac{1}{2}, \frac{1}{2}\right]$, then $\frac{s_1}{s_2}\xi \in \mathbb{R} \cup i \left[-\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right] \subset \mathbb{R} \cup i \left[-\frac{1}{2}, \frac{1}{2}\right]$ so that $\hat{\phi}(\xi) = \frac{s_1}{s_2}\hat{f}\left(\frac{s_1}{s_2}\xi\right) \geq 0$. It remains to estimate the resulting bounds on K. We compute

$$\frac{\sinh(s_1/2)}{s_1} \int_0^\infty \widehat{\phi}(x) x \tanh(\pi x) \, dx = \frac{\sinh(s_1/2)}{s_1} \frac{s_1}{s_2} \int_0^\infty \widehat{f}\left(\frac{s_1}{s_2}x\right) x \tanh(\pi x) \, dx$$
$$= \frac{\sinh(s_1/2)s_2}{s_1^2} \int_0^\infty \widehat{f}(y) y \tanh\left(\frac{s_2}{s_1}\pi y\right) \, dy$$
$$\leq \frac{\sinh(s_1/2)s_2^2}{s_1^3} \int_0^\infty \widehat{f}(y) y \tanh(\pi y) \, dy.$$

We then claim that

$$\frac{\sinh(s_1/2)s_2^2}{s_1^3} \le \frac{\sinh(s_2/2)}{s_2}$$

Indeed, this is equivalent to saying that $\sinh(x/2)/x^3$ increases from s_1 to s_2 and elementary computations show that the derivative of this function is non-negative provided that x is at least $6 \tanh(x/2)$, which is certainly true if $x \ge 6$.

We conclude that

$$\frac{\sinh(s_1/2)}{s_1|\phi(s_1)|} \int_0^\infty \widehat{\phi}(x)x \tanh(\pi x) \, dx \le \frac{\sinh(s_2/2)}{s_2|f(s_2)|} \int_0^\infty \widehat{f}(y)y \tanh(\pi y) \, dy$$

and hence that $K(s_1) \le K(s_2)$ upon taking the infimum over f .

When the systole is sufficiently small, any two shortest closed geodesics are disjoint, which implies a bound on the kissing number for topological reasons. More precisely, by the collar lemma [Bus10, Theorem 4.1.6], if $sys(M) \leq 2 \operatorname{arcsinh}(1)$, then $kiss(M) \leq 6(g-1)$ (and this is optimal).

In order to obtain a global upper bound on the kissing number in a given genus g from Theorem 7.1, it remains to prove an upper bound for K(s) when $s \in [2 \operatorname{arcsinh}(1), 6]$ and at $s = \operatorname{sys_bound}(g)$ where $\operatorname{sys_bound}(g)$ is the bound on the systole in genus g coming from the previous section. This is if $\operatorname{sys_bound}(g) > 6$, which happens as soon as $g \ge 7$. In lower genus, we only have to bound K(s) for s in $[2 \operatorname{arcsinh}(1), \operatorname{sys_bound}(g)]$.

7.2. Low genus. To deal with the interval $[2 \operatorname{arcsinh}(1), 6]$, we partition it into smaller subintervals, and then optimize to find a single function f for each subinterval. One subtlety is that to get an upper bound on K(s) when s belongs to an interval [a, b], we need an upper bound on

$$\frac{\sinh(s/2)}{-sf(s)}$$

for s in the same interval. Once we find a candidate upper bound numerically, we approximate $\sinh(s/2)$ and $e^{-s^2/2}$ by Taylor polynomials from the correct direction and then use Sturm's theorem to certify that the bound is valid (recall that $f(s) = p(s^2)e^{-s^2/2}$ for some polynomial p).

In other words, if we know that f(s) < 0 for every $s \in [a, b]$, then

$$\frac{\sinh(s/2)}{-sf(s)} \le B \quad \iff \quad \sinh(s/2) + Bs \, p(s^2) e^{-s^2/2} \le 0.$$

If $S_n(x) \ge \sinh(x/2)$ and $E_n(x) \le e^{-x^2/2}$ are polynomial approximations, then it suffices to check that

$$S_n(x) + Bx \, p(x^2)E_n(x) < 0,$$

on [a, b] (we are assuming that $p(x^2) < 0$ there, hence the reversed inequality for E_n). We can perform this verification by checking that $S_n(x) + Bx p(x^2)E_n(x)$ is negative at one point (using interval arithmetic) and has no zeros in [a, b] (using Sturm's theorem). For E_n , we simply take the odd degree Taylor polynomials of the exponential evaluated at $-x^2/2$ (the resulting series is alternating). For S_n , we use Taylor's theorem with the Lagrange form of the remainder to get

$$\sinh(x/2) = \sum_{j=0}^{n-1} \frac{(x/2)^{2j+1}}{(2j+1)!} + \cosh(\xi_x) \frac{(x/2)^{2n+1}}{(2n+1)!}$$

for some $\xi_x \in [0, x]$ and thus

$$\sinh(x/2) \le \sum_{j=0}^{n-1} \frac{(x/2)^{2j+1}}{(2j+1)!} + \cosh(b) \frac{(b/2)^{2n+1}}{(2n+1)!}$$

for every $x \in [a, b] \subset [0, b]$. We define $S_n(x)$ by the above formula except that we replace the last term involving b by a rational approximation from above, so that the computer can apply Sturm's theorem reliably.

The upper bounds on $\frac{1}{2}$ kiss that result from the above strategy are shown in Table 2 for genus 2 to 20 and verified in the ancillary file verify_kissing.ipynb. They improve upon the previous best bounds in every genus except g = 2, where the optimal bound is 12 for the number of unoriented systoles [Sch94]. We remark that these upper bounds depend in a very sensitive way on the upper bounds on the systole from Table 1, which are not as small as possible because we took precautions to make sure that they were rigorous. Consequently, the upper bounds in Table 2 could be decreased with more precision (especially those towards the end of the table).

7.3. Asymptotics. To prove an asymptotic upper bound the kissing number, we start with a proposition that bounds this quantity in terms of the systole.

Proposition 7.5. There exist some $s_0 > 0$ such that

$$K(s) < 2.922 \cdot \frac{2\sinh(s/2)}{s}$$

for every $s \ge s_0$. In particular, every closed hyperbolic surface M of genus $g \ge 2$ and systole $sys(M) \ge s_0$ satisfies

$$\operatorname{kiss}(M) < 2.922 \cdot \frac{2 \sinh(\operatorname{sys}(M)/2)}{\operatorname{sys}(M)} (g-1).$$

Remark 7.6. This improves upon [FBP22, Remark 4.4], where we obtained the same inequality but with the constant 63.71 instead of 2.922.

Proof. Recall that

$$K(s) = 2\sqrt{2\pi} \frac{2\sinh(s/2)}{s} \inf_{f} \left\{ \frac{1}{-f(s)} \int_{0}^{\infty} \widehat{f}(x)x \tanh(x) \, dx \right\}$$

where the infimum is over admissible functions f such that $f(x) \leq 0$ if $x \geq s$, f(s) < 0, and $\widehat{f}(\xi) \geq 0$ if $\xi \in \mathbb{R} \cup i[-\frac{1}{2}, \frac{1}{2}]$.

genus	lower bound	LP bound	previous upper bound
2	12 [Jen84]	14^{\dagger}	12 [Sch94]
3	24 [Sch93]	34	126 [MRT14]
4	36 [Sch93]	62	244 [FBP22]
5	48 [Sch93]	97	383 [FBP22]
6	$39 \left[\mathrm{Ham}01 \right]$	138	547 [FBP22]
7	126 [Vog03]	185	736 [FBP22]
8		240	950 [FBP22]
9	70 [Sch93]	299	1186 [FBP22]
10		364	1446 [FBP22]
11	120 [Sch93]	434	1728 [FBP22]
12		510	2032 [FBP22]
13	144 [FBR22]	591	2358 [FBP22]
14	364 [Vog03]	677	2706 [FBP22]
15	168 [FBR22]	771	3074 [FBP22]
16	180 [FBR22]	868	3464 [FBP22]
17	336 [Vog03]	970	3874 [FBP22]
18	204 [FBR22]	1083	4305 [FBP22]
19	216 [FBR22]	1209	4756 [FBP22]
20	228 [FBR22]	1333	5227 [FBP22]

TABLE 2. Bounds on the maximum of $\frac{1}{2}$ kiss. The dagger indicates where the linear programming bound fails to beat the previous best upper bound.

To prove the desired inequality, we use the same functions as for the asymptotic systole bound but choose the parameters differently. That is, we take f such that

$$\widehat{f}(x) = (c - 1 + b^2 x^2) e^{-b^2 x^2/2} \varphi_{\alpha}(Rx) = (c - 1 + b^2 x^2) e^{-b^2 x^2/2} \frac{J_{\alpha}(Rx/2)^2}{(Rx)^{2\alpha}}$$

for some parameters α , b, c and R but now set $c = 1 + b^2/4$ so that $\widehat{f}(i/2) = 0$.

Recall from the proof of Lemma 6.7 that for any $\kappa_0 \in \mathbb{R}$ we have that $R^{2\alpha+1}f(R+b\kappa_0)$ tends to

$$-L_1 = -\frac{4b^{2\alpha}\cos\left(\frac{(2\alpha+1)\pi}{2}\right)}{\pi} \int_0^\infty \frac{1-\cos(x)}{x^{2\alpha+1}} dx \int_{\kappa_0}^\infty (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} dx,$$

as $R \to \infty$. We also saw in the proof of Theorem 6.4 that

$$L_2 = \lim_{R \to \infty} R^{2\alpha + 1} \int_0^\infty \widehat{f}(x) x \tanh(\pi x) \, dx = \frac{2b^2}{\pi} \int_0^\infty \frac{\tanh(\pi x)}{x^{2\alpha}} (x^2 + 1/4) e^{-b^2 x^2/2} \, dx.$$

Our goal is thus to minimize the ratio L_2/L_1 over the parameters such that $-L_1 < 0$ and such that f is non-positive from $s = R + b\kappa_0$ onwards. At $\alpha = 0.592$, b = 0.981 and $\kappa_0 = 0.061$, we obtain $2\sqrt{2\pi}L_2/L_1 < 2.922$, which implies that

$$2\sqrt{2\pi}\frac{1}{-f(R+b\kappa_0)}\int_0^\infty \widehat{f}(x)x\tanh(x)\,dx < 2.922$$

if R is large enough. To prove that $2\sqrt{2\pi}L_2/L_1 < 2.922$, we observe that

$$\int_0^\infty \frac{\tanh(\pi x)}{x^{2\alpha}} (x^2 + 1/4) e^{-b^2 x^2/2} \, dx$$

is bounded above by

$$\pi \left(A^2 + \frac{1}{4}\right) \frac{A^{2(1-\alpha)}}{2(1-\alpha)} + \int_A^B \frac{\tanh(\pi x)}{x^{2\alpha}} (x^2 + 1/4) e^{-b^2 x^2/2} \, dx + \frac{1}{B^{2\alpha}} \int_B^\infty x^3 e^{-b^2 x^2/2} \, dx$$

whenever 0 < A < 2 < B and estimate each term using interval arithmetic with $A = 1/10^4$ and $B = 10^4$, noting that the last term can be rewritten as

$$\frac{1}{B^{2\alpha}} \int_B^\infty x^3 e^{-b^2 x^2/2} \, dx = \frac{1}{b^4 B^{2\alpha}} \int_{bB}^\infty y^3 e^{-y^2/2} \, dy = \frac{(b^2 B^2 + 2)}{b^4 B^{2\alpha}} e^{-b^2 B^2/2}.$$

For L_1 , we need a lower bound. We have

$$\int_0^\infty \frac{1 - \cos(x)}{x^{2\alpha + 1}} dx \ge \int_A^B \frac{1 - \cos(x)}{x^{2\alpha + 1}} dx$$

since the integrand is non-negative. We then observe that

$$-\int_{\kappa_0}^{\infty} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx \ge -\int_{\kappa_0}^{B} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx,$$

since $B > \sqrt{c}$. We split this last integral at $\kappa_0 + A$ and estimate

$$-\int_{\kappa_0}^{\kappa_0+A} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} \, dx \ge -A^{2\alpha} \int_{\kappa_0}^{\kappa_0+A} (c-x^2) e^{-x^2/2} \, dx$$

Putting all the estimates together, we obtain the certified upper bound

$$2\sqrt{2\pi}\frac{L_2}{L_1} \le 2.92190512955185 < 2.922.$$

The last thing to check is that when R is large enough, we have $f(x) \leq 0$ for every $x \geq s = R + b\kappa_0$. The proof of this fact is similar as for the systole bound, and is deferred to the next lemma.

The resulting upper bound on kiss(M) when $sys(M) \ge s_0$ then follows from Theorem 7.1.

We now prove a small lemma which verifies that the function f used above satisfies the hypotheses of Theorem 7.1 for $s = R + b\kappa_0$ whenever R is large enough. This is similar to Lemma 6.9 and Corollary 6.10.

Lemma 7.7. Let f be as in Proposition 7.5 with $\alpha = 0.592$, b = 0.981, $c = 1 + b^2/4$, and $\kappa_0 = 0.061$. Then there exists some $R_0 > 0$ such that $f(x) \leq 0$ for every $x \geq R + b\kappa_0$ and every $R \geq R_0$.

Proof. We begin by proving the pointwise result that for every $\kappa \geq \kappa_0$, we have that $f(R + b\kappa) \leq 0$ if R is sufficiently large. Since f is the convolution of a non-negative function supported in [-R, R] and a function which is non-positive outside $[-b\sqrt{c}, b\sqrt{c}]$, it is obviously non-positive at points x with $|x| \geq R + b\sqrt{c}$. In other words, the result is obvious (and holds for every R > 0) if $\kappa \geq \sqrt{c}$.

For $\kappa \in [\kappa_0, \sqrt{c}]$, we use the fact that $R^{2\alpha+1}f(R+b\kappa)$ converges to a positive multiple of

$$I(\kappa) = \int_{\kappa}^{\infty} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} dx$$

as $R \to \infty$, so it suffices to check that $I(\kappa) < 0$ for every $\kappa \in [\kappa_0, \sqrt{c}]$. This is similar to the statement of Lemma 6.9 but is easier to prove because $I(\kappa_0)$ is not close to 0, so coarse bounds suffice. We write

$$I(\kappa) = \int_{\kappa}^{\sqrt{c}} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} dx + \int_{\sqrt{c}}^{\infty} (x-\kappa)^{2\alpha} (c-x^2) e^{-x^2/2} dx$$
$$\leq \int_{\kappa_0}^{\sqrt{c}} (x-\kappa_0)^{2\alpha} (c-x^2) e^{-x^2/2} dx + \int_{\sqrt{c}}^{\infty} (x-\sqrt{c})^{2\alpha} (c-x^2) e^{-x^2/2} dx$$

then split the first integral at $\kappa_0 + 1/100$ and the second one at 100 to get that

$$I(\kappa) \le -0.276593809735452 < 0$$

for every $\kappa \in [\kappa_0, \sqrt{c}]$.

From this pointwise result and the same continuity and compactness argument as in Corollary 6.10, we obtain that there exists some $R_0 > 0$ such that $f(R + b\kappa) \leq 0$ for every $\kappa \geq \kappa_0$ and every $R \geq R_0$. In other words, we have that $f(x) \leq 0$ for every $x \geq R + b\kappa$ and every $R \geq R_0$.

We then combine the previous proposition with our asymptotic bound for the systole to obtain the following bound on kissing numbers that only depends on the genus.

Theorem 7.8. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

kiss
$$(M) < \frac{4.873 \cdot g^2}{\log(g) + 1.2045}.$$

Remark 7.9. For some towers of principal congruence covers of arithmetic surfaces, the kissing number grows at least like $g^{4/3-\varepsilon}$ for any $\varepsilon > 0$ [SS97]. This was recently generalized to higher dimensional hyperbolic manifolds in [DFM22] with g replaced by the volume and an exponent that depends on the dimension.

Proof of Theorem 7.8. Recall that Theorem 6.4 states that

$$\operatorname{sys}(M) \le 2\log(g) + 2.409$$

if g is large enough. Let s_0 be as in Proposition 7.5. If g is sufficiently large, then $2\log(g) + 2.409 \ge s_0$ and we get

$$\begin{split} K(2\log(g)+2.409) &\leq 2.922 \, \frac{2\sinh((2\log(g)+2.409)/2)}{2\log(g)+2.409} \\ &< 1.461 \, \frac{e^{1.2045}g}{\log(g)+1.2045} \\ &< \frac{4.873 \cdot g}{\log(g)+1.2045} \end{split}$$

where we used the fact that $2\sinh(x/2) < e^{x/2}$ for every real number x.

By Theorem 7.1, we have

$$kiss(M) \le K(sys(M))(g-1).$$

Furthermore, Lemma 7.4 says that K is non-decreasing on $[6, \infty)$. We thus get that

kiss
$$(M) \le K(2\log(g) + 2.409)(g-1) < \frac{4.873 \cdot g^2}{\log(g) + 1.2045}$$

provided $sys(M) \ge 6$ and $2\log(g) + 2.409 \ge \max\{6, s_0\}$, which holds whenever g is large enough.

When the systole is at most $2 \operatorname{arcsinh}(1)$, we noted previously that $\operatorname{kiss}(M) \leq 6(g-1)$ by the collar lemma, and this quantity is smaller than the stated bound when g is large enough (in fact, this is true for all $g \geq 2$).

The only interval left to cover is $[2 \operatorname{arcsinh}(1), 6]$. By the calculations used to produce Table 2, the function K is bounded on that interval. One can also prove this using a single function f defined by

$$\widehat{f}(x) = (x^2 + 1/4)e^{-x^2/2}\varphi_{\alpha}(Rx)$$

with $R = 2 \operatorname{arcsinh}(1) - \sqrt{5/4} > 0$ and any $\alpha \in (0, 1)$, because it satisfies all the hypotheses of Theorem 7.1 for every $s \ge 2 \operatorname{arcsinh}(1)$. The resulting linear upper bound on kiss(M) when sys $(M) \in [2 \operatorname{arcsinh}(1), 6]$ is eventually smaller than $\frac{4.873 \cdot g^2}{\log(g) + 1.2045}$ when g is large enough.

8. First eigenvalue

8.1. The criterion. The criterion for bounding the first positive eigenvalue $\lambda_1(M)$ of the Laplacian on M goes as follows.

Theorem 8.1. Let $g \ge 2$. Suppose that f is a non-constant admissible function for which there exists an L > 0 such that

• $f(x) \ge 0$ for all $x \in \mathbb{R}$;

•
$$\widehat{f}\left(\sqrt{\lambda-\frac{1}{4}}\right) \leq 0$$
 whenever $\lambda \geq L$;

• $\widehat{f}(i/2) \le 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) dx;$

Then $\lambda_1(M) \leq L$ for every hyperbolic surface M of genus g.

Proof. Suppose that M is a hyperbolic surface with $\lambda_1(M) > L$. By continuity, the same inequality holds for every surface N in some neighborhood U of M in moduli space. Let $r_j(N) \in \mathbb{C}$ be such that $r_j(N)^2 = \lambda_j(N) - \frac{1}{4}$. The Selberg trace formula yields

$$\begin{split} \sum_{j=0}^{\infty} \widehat{f}(r_j(N)) &\leq \widehat{f}(i/2) \\ &\leq 2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx \\ &\leq 2(g-1) \int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(N)} \frac{\Lambda(\gamma) f(\ell(\gamma))}{2 \sinh(\ell(\gamma)/2)} \\ &= \sum_{j=0}^{\infty} \widehat{f}(r_j(N)) \end{split}$$

from which we conclude that $\hat{f}(r_j(N)) = 0$ for every $j \ge 1$ and every $N \in U$. As in the proof of Theorem 6.1, this leads to a contradiction since the zero set of \hat{f} is discrete and there are no non-trivial isospectral deformations of a hyperbolic surface. We conclude that $\lambda_1(M) \le L$ for every M.

Remark 8.2. If the inequality in the third bullet point is strict, then the conclusion can be strengthened to a strict inequality and this is easier to prove.

8.2. Low genus. The upper bounds on $\lambda_1(M)$ resulting from Theorem 8.1 and numerical optimization are presented in Table 3 for $2 \leq g \leq 20$ and their verification is done in the file verify_lambda.ipynb. Our bounds are smaller than the previous best upper bounds in every genus except 2, 3, 4, and 6 where bounds from [KMP21], [Bon22] or [YY80] are better.

genus	lower bound	LP bound	previous upper bound
2	3.838887 [SU13]	4.625307^{\dagger}	3.838898 [KMP21, Bon22]
3	2.6779 [Coo18]	2.816427^{\dagger}	2.678483 [KMP21, Bon22]
4	1.91556 [Coo18]	2.173806^{\dagger}	2.000000 [YY80]
5	0.728167 (§10.4)	1.836766	1.852651 [KMP21]
6	0.486360 (§10.4)	1.625596^{\dagger}	1.600000 [YY80]
7	0.691340 (§10.4)	1.480008	1.513268 [KMP21]
8	0.25 [BS07] + [BBD88]	1.372804	1.406905 [KMP21]
9	0.25 [BS07] + [BBD88]	1.289024	1.323482 [KMP21]
10	0.25 [BS07] + [BBD88]	1.222189	1.256022 [KMP21]
11	0.25 [BS07]+[BBD88]	1.168169	1.200153 [KMP21]
12	0.25 [BS07]+[BBD88]	1.122327	1.152986 [KMP21]
13	0.25 [BS07]+[BBD88]	1.083260	1.112535 [KMP21]
14	0.287470 (§10.4)	1.049217	1.077385 [KMP21]
15	0.25 [BS07] + [BBD88]	1.018005	1.046501 [KMP21]
16	0.25 [BS07]+[BBD88]	0.991735	1.019105 [KMP21]
17	0.403200 (§10.4)	0.968260	0.994601 [KMP21]
18	0.25 [BS07]+[BBD88]	0.947180	0.972525 [KMP21]
19	0.25 [BS07]+[BBD88]	0.928091	0.952510 [KMP21]
20	0.25 [BS07]+[BBD88]	0.911390	0.934260 [KMP21]

TABLE 3. Bounds on the supremum of λ_1 . The daggers indicate where the linear programming bounds fail to beat the previous best upper bound.

Some comments on the examples we used for the lower bounds are in order:

- Contrary to the other invariants considered in this paper, it is not known if the supremum of λ_1 is attained. For instance, we do not know if the entries equal to 1/4 in the table are attained (see below).
- In genus 2 and 3, the upper bounds from [KMP21, Bon22] are tantalizingly close to the value of λ_1 at the Bolza surface and the Klein quartic approximated numerically in [SU13] and [Coo18] respectively, so these surfaces are the conjectured maximizers in these genera. The authors of [KMP21] reproduced Cook's numerical calculations with more precision, arriving at the value 2.6779 instead of 2.6767 for the Klein

quartic. The surface in genus 4 is Bring's curve. Note that the values in genus 3 and 4 are based on finite element methods and are not rigorous. The value in genus 2 is obtained using the trace formula and can be made rigorous according to Strohmaier and Uski.

- In genus 5 to 7, 14, and 17, we apply linear programming to some of the surfaces listed in Table 1 to obtain lower bounds on their first eigenvalue based on their systole (see Section 10.4). The true value of λ_1 for these surfaces is certainly larger than the estimates we give since we discard all the geometric terms and the contribution of higher eigenvalues in the Selberg trace formula, while the test functions we use have only finitely many zeros. For instance, preliminary numerical calculations by Master's student Mathieu Pineault indicate that the Fricke–Macbeath curve in genus 7 has $\lambda_1 \approx 1.239$.
- If $X \to Y$ is a finite-sheeted covering of hyperbolic orbifolds, then $\lambda_1(Y) \ge \lambda_1(X)$ since any eigenfunction on Y lifts to an eigenfunction on X with the same eigenvalue. We will use this in the next bullet point.
- For the entries equal to 1/4 in the table, we use the fact that Selberg's conjecture is known to hold for the congruence subgroups $\Gamma_1(N)$ of square-free level N < 857[BS07]. Since $\Gamma_1(N) < \Gamma_0(N)$ as a finite-index subgroup, the conjecture also holds for $\Gamma_0(N)$ for the same levels. If $\Gamma_j(N)$ is torsion-free, then $X_j(N) := \mathbb{H}^2/\Gamma_j(N)$ (j = 0, 1) has no cone points and we can join its cusps in pairs to create thin handles following [BBD88]. By the results in that paper, the spectrum of the plumbed surface will be close to that of $X_j(N)$ (which has a discrete spectrum and a continuous spectrum equal to $[1/4, \infty)$ like all cusped surfaces). Since the spectral gap of $X_j(N)$ is 1/4 for square-free N < 857 [BS07], the plumbed surfaces thus have λ_1 as close to 1/4 as we wish in these cases. The genus of the plumbed surfaces is equal to the genus of $X_j(N)$ plus half its number of cusps. Table 4 shows which congruence group we use for each genus concerned. The fact that these groups are indeed torsion free and that their signatures are as listed can be found in [Miy06, Section 4.2]. This information about congruence groups is also implemented in Sage.
- There are other well-known ways of proving lower bounds on λ_1 . The first of these is Cheeger's inequality $\lambda_1 \geq h^2/4$ where h is the Cheeger constant [Che70]. In high genus, this cannot be used to prove a lower bound of more than $\frac{1}{\pi^2} \approx 0.1013...$ on λ_1 [BCP22]. However, it is not clear what bounds this approach might give in low genus as there are no explicit calculations of Cheeger constants for closed surfaces yet [AM99, Ben15, BLT21]. The second approach is to use the Jacquet– Langlands correspondence [JL70], which allows one to derive lower bounds on the first eigenvalue of certain compact arithmetic surfaces from lower bounds on the first discrete eigenvalue of corresponding congruence covers of the modular curve (see [Ber16, Example 8.27] and [Hej85] for concrete examples). We are not aware of any examples where both of the following hold: a better lower bound than 1/4 is known for the cusped surface (see [Hux85, BSV06] for examples) and the Jacquet– Langlands correspondence gives rise to a closed surface of genus at most 20 without cone points.

Γ	(g,n)	$g_{ m plumbed}$
$\Gamma_1(13)$	(2, 12)	8
$\Gamma_1(15)$	(1, 16)	9
$\Gamma_0(107)$	(9, 2)	10
$\Gamma_0(87)$	(9, 4)	11
$\Gamma_0(86)$	(10, 4)	12
$\Gamma_1(17)$	(5, 16)	13
$\Gamma_0(78)$	(11, 8)	15
$\Gamma_1(19)$	(7, 18)	16
$\Gamma_0(134)$	(16, 4)	18
$\Gamma_0(102)$	(15, 8)	19
$\Gamma_0(227)$	(19, 2)	20

TABLE 4. Some congruence subgroups of the modular group, their signatures, and the genus of their plumbing.

8.3. Asymptotics. A theorem of Cheng [Che75, Theorem 2.1] states that

(8.1)
$$\lambda_1(M) \le \lambda_0(D_{\operatorname{diam}(M)/2})$$

for any closed hyperbolic surface M, where D_R is a hyperbolic disk of radius R, $\lambda_0(\Omega)$ is the smallest Dirichlet eigenvalue of Ω , and diam(M) is the diameter of M. From this, Cheng deduces [Che75, Corollary 2.3] the more explicit bound

$$\lambda_1(M) \le \frac{1}{4} + \left(\frac{4\pi}{\operatorname{diam}(M)}\right)^2.$$

However, this can be improved using an inequality of Gage [Gag80, Theorem 5.2(a)] on the smallest eigenvalue of hyperbolic disks, which states that

(8.2)
$$\lambda_0(D_R) \le \frac{1}{4} + \frac{\pi^2}{R^2} - \frac{1}{4\sinh^2(R)}$$

If we ignore the last term (of smaller order), we obtain the improved inequality

$$\lambda_1(M) \le \frac{1}{4} + \left(\frac{2\pi}{\operatorname{diam}(M)}\right)^2$$

In turn, the best known lower bound on the diameter is Bavard's bound

diam
$$(M) \ge \operatorname{arccosh}\left(\frac{1}{\sqrt{3}\tan(\pi/(12g-6))}\right)$$

where g is the genus of M [Bav96]. Since the $\tan(x) \sim x$ as $x \to 0$ and $\operatorname{arccosh}(x)$ is asymptotic to $\log(x)$ as $x \to \infty$, Bavard's inequality has the same asymptotic behaviour as the more elementary inequalities

$$\operatorname{diam}(M) \ge 2\operatorname{arcsinh}(\sqrt{g-1}) \ge \log(g-1)$$

coming from area considerations, which result in

$$\lambda_1(M) \le \frac{1}{4} + \left(\frac{2\pi}{\log(g-1)}\right)^2.$$

We will improve upon this by another factor of 4.

Theorem 8.3. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$\lambda_1(M) < \frac{1}{4} + \left(\frac{\pi}{\log(g) + 0.7436}\right)^2.$$

Remark 8.4. An inequality of Savo [Sav09, Theorem 5.6] states that

(8.3)
$$\lambda_0(D_R) \ge \frac{1}{4} + \frac{\pi^2}{R^2} - \frac{4\pi^2}{R^3}$$

for every R > 0. It follows that the leading terms in Gage's upper bound (8.2) cannot be improved. Moreover, there exist sequences of hyperbolic surfaces M_g of genus g whose diameter is asymptotic to $\log(g)$ [BCP21]. It follows that no multiplicative improvement as in Theorem 8.3 could be obtained from Cheng's inequality (8.1).

Remark 8.5. As stated in the introduction, it is still unknown if there exist surfaces with $\lambda_1(M) \geq 1/4$ in every genus. However, it was proved recently that for any $\varepsilon > 0$, there exist surfaces with $\lambda_1(M) > 1/4 - \varepsilon$ in large enough genus [HM21].

The proof of Theorem 8.3 will require the following technical lemma whose proof is postponed until after the proof of the theorem.

Lemma 8.6. We have

$$\lim_{R \to \infty} R^4 \int_0^\infty \frac{\sin^2(\pi Rx)}{1 - (Rx)^2} \frac{x \tanh(\pi x)}{(Rx)^2} \, dx = \frac{\pi^2}{2} \int_0^\infty \frac{\sinh(x) \cosh(x) - x}{x^3 \cosh^2(x)} \, dx > 4.20718.$$

We use this to prove the theorem.

Proof of Theorem 8.3. The idea of the proof is to apply Theorem 8.1 with functions of the form $f_R(x) = f(x/R)/R$ for some fixed non-negative admissible function f such that \hat{f} is non-positive on $[1, \infty)$, so that $\hat{f}_R(x) = \hat{f}(Rx)$ is non-positive on $[1/R, \infty)$. Then R = R(g) must be chosen as large as possible such that the inequality

(8.4)
$$\widehat{f}(Ri/2) \le 2(g-1) \int_0^\infty \widehat{f}(Rx) x \tanh(\pi x) \, dx$$

remains valid.

We choose

$$f(x) = \sqrt{\frac{\pi}{8}} (2\pi - |x| + \sin|x|) \chi_{[-2\pi, 2\pi]}(x)$$

whose Fourier transform is equal to

$$\widehat{f}(x) = \frac{\sin^2(\pi x)}{x^2(1-x^2)}.$$

Note that f is non-negative on \mathbb{R} and \hat{f} is non-positive on $[1,\infty)$, as required. We also have that f is admissible since \hat{f} is entire and is $O(|x|^{-4})$ on any horizontal strip of finite height.

For a given genus $g \ge 2$, we want to find an R = R(g) > 0 such that inequality (8.4) holds. We first compute

$$\widehat{f}(Ri/2) = \frac{4\sinh^2(\pi R/2)}{R^2(1+R^2/4)} < \frac{4e^{\pi R}}{R^4}.$$

For the integral term, we have

$$\lim_{R \to \infty} R^4 \int_0^\infty \widehat{f}(Rx) x \tanh(\pi x) \, dx = \frac{\pi^2}{2} \int_0^\infty \frac{\sinh(x) \cosh(x) - x}{x^3 \cosh^2(x)} dx$$

by Lemma 8.6. If c is any positive number strictly smaller than the right-hand side (such as 4.2071), then we have

(8.5)
$$\frac{\widehat{f}(Ri/2)}{\int_0^\infty \widehat{f}(Rx)x \tanh(\pi x) \, dx} < \frac{4}{c} e^{\pi R}$$

provided that R is large enough. If g is sufficiently large, then we can take $R = \frac{\log(g-1) + \log(c/2)}{\pi}$ so that the right-hand side of (8.5) becomes equal to 2(g-1). Then f_R satisfies the hypotheses of Theorem 8.1, which proves the upper bound $\lambda_1(M) \leq L$ for every closed hyperbolic surface M of large enough genus g, where L is such that

$$\sqrt{L-1/4} = 1/R$$

the point after which $\widehat{f_R}$ stays non-positive. This gives

$$L = \frac{1}{4} + \frac{1}{R^2} = \frac{1}{4} + \frac{\pi^2}{(\log(g-1) + \log(c/2))^2}.$$

Since $\log(4.20718/2) > 0.7436$ and $\log(g) - \log(g-1)$ tends to zero as $g \to \infty$, the inequality

$$\lambda_1(M) < \frac{1}{4} + \frac{\pi^2}{(\log(g) + 0.7436)^2}$$

holds for all closed hyperbolic surfaces of sufficiently large genus.

Remark 8.7. The choice of f in the above proof is not at all random. It is proportional to the function $\psi_{1/2}(j_{1/2}x)$ from Section 5. It can be shown that this function is optimal for the strategy we use. Indeed, the problem amounts to minimizing the growth of $\hat{f}(Ri/2)$ among functions such that $\int_0^{\infty} \hat{f}(Rx)x \tanh(\pi x) dx$ is eventually positive (so that equation (8.4) has any chance of being satisfied). By a change of variable and an application of the dominated convergence theorem (see the proof of Proposition 9.3 below), we get that a necessary condition for this eventual positivity is that the second moment $\int_0^{\infty} \hat{f}(x)x^2 dx$ is non-negative. Moreover, by the Paley–Wiener theorem, the growth of $\hat{f}(Ri/2)$ is controlled by the support of f. The question thus boils down to minimizing the support of f among non-negative even functions whose Fourier transform is non-positive outside [-1, 1] and has a non-negative second moment. According to [GIT20, Remark 1.2] (with d = s = 1, m = 0, and f and \hat{f} interchanged), the optimal function for this problem is the one we used.

We now prove the technical lemma.

Proof of Lemma 8.6. We start by making the change of variable x = Ry to get

$$R^{2} \int_{0}^{\infty} \frac{\sin^{2}(\pi Rx)}{1 - (Rx)^{2}} \frac{\tanh(\pi x)}{x} \, dx = \pi R \int_{0}^{\infty} \frac{\sin^{2}(\pi y)}{1 - y^{2}} \frac{\tanh(\pi y/R)}{\pi y/R} \, dy.$$

Let $F(y) = \frac{\sin^2(\pi y)}{y}$, $G_R(y) = R \frac{\tanh(\pi y/R)}{\pi y/R}$, and $H_R(u) = G_R(u-1) - G_R(u+1)$. Observe that F is odd, G_R is even, and H_R is odd. For every b > 1, we have

$$2R \int_0^b \frac{\sin^2(\pi y)}{1 - y^2} \frac{\tanh(\pi y/R)}{\pi y/R} \, dy = \int_1^{b+1} F(u) G_R(u-1) \, du - \int_1^{1-b} F(v) G_R(v-1) \, dv$$
$$= \int_{1-b}^{b+1} F(u) G_R(u-1) \, du$$

by the changes of variable u = 1 + y and v = 1 - y. We can rewrite this as

$$\int_{b-1}^{b+1} F(u)G_R(u-1)\,du + \int_{1-b}^{b-1} F(u)G_R(u-1)\,du$$

and by breaking up the second integral at u = 0 and using the symmetries of F and G_R , we obtain

$$\int_{b-1}^{b+1} F(u)G_R(u-1)\,du + \int_0^{b-1} F(u)H_R(u)du.$$

For every fixed R, we have that $G_R(y)$ is bounded between 0 and R so that

$$\left| \int_{b-1}^{b+1} F(u) G_R(u-1) \, du \right| \le R \int_{b-1}^{b+1} |F(u)| \, du$$
$$\le \int_{b-1}^{b+1} \frac{1}{u} \, du$$
$$= \log\left(\frac{b+1}{b-1}\right)$$

tends to zero as $b \to \infty$. We thus have

$$2R \int_0^\infty \frac{\sin^2(\pi y)}{1 - y^2} \frac{\tanh(\pi y/R)}{\pi y/R} \, dy = \int_0^\infty F(u) H_R(u) \, du.$$

Since $2uF(u) = 2\sin^2(\pi u) = 1 - \cos(2\pi u)$, we obtain

$$2\int_{0}^{\infty} F(u)H_{R}(u) \, du = \int_{0}^{\infty} \frac{H_{R}(u)}{u} du - \int_{0}^{\infty} \cos(2\pi u) \frac{H_{R}(u)}{u} du$$
$$= \int_{0}^{\infty} \frac{H_{R}(Rx/\pi)}{x} dx - \int_{0}^{\infty} \cos(2Rx) \frac{H_{R}(Rx/\pi)}{x} dx$$

by the change of variable $u = Rx/\pi$. We will show that the second term tends to zero as $R \to \infty$ and come back to the first term afterwards. If $\phi_R(x) := H_R(Rx/\pi)/x$ did not depend on R, this would follow directly from the Riemann–Lebesgue lemma. However, the proof of the Riemann–Lebesgue lemma still implies the desired result with some work. Since f_R is even, we have

$$2\int_0^\infty \cos(2Rx)\phi_R(x)dx = \int_{-\infty}^\infty \cos(2Rx)f_R(x)dx$$

By the change of variable $x = y + \frac{\pi}{2R}$, we get

$$\int_{-\infty}^{\infty} \cos(2Rx)\phi_R(x)dx = -\int_{-\infty}^{\infty} \cos(2Ry)\phi_R\left(y + \frac{\pi}{2R}\right)dy$$

so that

$$\int_{-\infty}^{\infty} \cos(2Rx)\phi_R(x)dx = \frac{1}{2}\int_{-\infty}^{\infty} \cos(2Rx)\left(\phi_R(x) - \phi_R\left(x + \frac{\pi}{2R}\right)\right)dx$$

It then suffices to check that $\phi_R(x) - \phi_R\left(x + \frac{\pi}{2R}\right)$ tends to zero as $R \to \infty$ and is bounded by an integrable function in absolute value. We have

$$\phi_R(x) = \frac{H_R(Rx/\pi)}{x} = \frac{G_R(Rx/\pi - 1) - G_R(Rx/\pi + 1)}{x}$$
$$= \frac{R}{x} \left(\frac{\tanh(x - \pi/R)}{x - \pi/R} - \frac{\tanh(x + \pi/R)}{x + \pi/R} \right)$$

Let $h(x) = \tanh(x)/x$ for $x \neq 0$ and h(0) = 1. Then for every $x \in \mathbb{R}$ we have

$$\lim_{R \to \infty} \frac{h(x - \pi/R) - h(x + \pi/R)}{-2\pi/R} = h'(x)$$

by definition of the derivative and similarly

$$\lim_{R \to \infty} \frac{h(x - \frac{\pi}{2R}) - h(x + \frac{3\pi}{2R})}{-2\pi/R} = h'(x)$$

by the mean value theorem and the continuity of h'. It follows that $\phi_R(x)$ and $\phi_R\left(x + \frac{\pi}{2R}\right)$ both converge to $-2\pi h'(x)/x$ as $R \to \infty$ and hence their difference tends to zero, at least for every $x \neq 0$.

It remains to show that $\phi_R(x)$ and $\phi_R\left(x + \frac{\pi}{2R}\right)$ are bounded above by integrable functions that do not depend on R. The Maclaurin series of h is $1 - \frac{x^2}{3} + O(x^4)$ so that

$$h(x - \pi/R) - h(x + \pi/R) = h(\pi/R - x) - h(\pi/R + x)$$

= $\frac{(x + \pi/R)^2}{3} - \frac{(x - \pi/R)^2}{3} + \frac{h''(\zeta)}{2}(-2x)^2$
= $\frac{4\pi x}{3R} + 2h''(\zeta)x^2$

for some $\zeta \in (\pi/R - x, \pi/R + x)$ by Taylor's theorem. Since all the derivatives of h are bounded on the real line, the error term is at most Ax^2 for some constant A > 0 that does not depend on R. This implies that

$$\frac{R}{x}(h(x-\pi/R) - h(x+\pi/R)) \le B$$

for some constant B > 0 whenever $x \in [0, 2\pi/R]$.

We now estimate the function away from the origin. For every $x \in \mathbb{R}$, there exists some $x^* \in (x - \pi/R, x + \pi/R)$ such that

$$h(x - \pi/R) - h(x + \pi/R) = -(2\pi/R)h'(x^*)$$

by the mean value theorem. We compute

$$-h'(x) = \frac{\tanh(x) - \frac{x}{\cosh^2(x)}}{x^2} \le \frac{\tanh(x)}{x^2} \le \frac{1}{x^2}$$

whenever x > 0 so that $-h'(x^*) \le 1/(x - \pi/R)^2$ and hence

$$\frac{R}{x}(h(x-\pi/R) - h(x+\pi/R)) \le \frac{2\pi}{x(x-\pi/R)^2}$$

provided that $x > \pi/R$. Also note that

$$\frac{1}{x - \pi/R} \le \frac{2}{x}$$

if $x \ge 2\pi/R$, which leads to the uniform estimate

$$\frac{R}{x}(h(x - \pi/R) - h(x + \pi/R)) \le \frac{8\pi}{x^3}$$

whenever $x \ge 2\pi/R$. This is only useful if $x \ge 1$, so we need a better bound on $[2\pi/R, 1]$. We can write

$$-h'(x) = \frac{\sinh(x)\cosh(x) - x}{(x\cosh(x))^2} \le \frac{\sinh(x)\cosh(x) - x}{x^2} = \frac{\frac{\sinh(2x)}{2} - x}{x^2}$$

By Taylor's theorem, there exists some C > 0 such that

$$\frac{\sinh(2x)}{2} \le x + \frac{2}{3}x^3 + Cx^5$$

whenever $x \in [0, 2]$. Therefore, if $x \in [2\pi/R, 1]$ and $R > \pi$, then

$$x^* \in (x - \pi/R, x + \pi/R) \subset (\pi/R, 1 + \pi/R) \subset [0, 2]$$

so that

$$-h'(x^*) \le \frac{2}{3}x^* + C \cdot (x^*)^3.$$

This gives the estimate

$$\frac{R}{x}(h(x-\pi/R) - h(x+\pi/R)) = -\frac{2\pi h'(x^*)}{x} \le 2\pi \frac{\frac{2}{3}(x+\pi/R) + C \cdot (x+\pi/R)^3}{x}$$

for every $x \in [2\pi/R, 1]$ provided that $R > \pi$. It is easy to check that $\frac{x+\pi/R}{x}$ is decreasing on that interval, hence is bounded by 3/2 while $\frac{(x+\pi/R)^3}{x}$ is increasing and hence bounded by $(1 + \pi/R)^3 \le 8$.

Putting these estimates together, we get that there exists a constant D > 0 such that $\phi_R(x) = \frac{R}{x}(h(x - \pi/R) - h(x + \pi/R))$ is bounded above by

$$\mu(x) = \begin{cases} D & \text{if } x \in [-1, 1] \\ \frac{8\pi}{|x|^3} & \text{if } |x| > 1 \end{cases}$$

for every $x \in \mathbb{R}$ provided that $R > 2\pi$. Note that μ is integrable. Similarly,

$$\phi_R\left(x+\frac{\pi}{2R}\right) \le \mu\left(x+\frac{\pi}{2R}\right) \le \nu(x)$$

where

$$\nu(x) = \begin{cases} \mu(x) & \text{if } x \ge -1\\ \mu(x+1/4) & \text{if } x < -1 \end{cases}$$

is still integrable. We can therefore apply the dominated convergence theorem to conclude that - - -

$$\lim_{R \to \infty} \int_{-\infty}^{\infty} \cos(2Rx) \left(\phi_R(x) - \phi_R\left(x + \frac{\pi}{2R}\right) \right) dx = 0$$

and hence

$$\lim_{R \to \infty} \int_0^\infty \cos(2Rx) \frac{H_R(Rx/\pi)}{x} dx = 0,$$

as claimed.

Returning to the original problem, we have

$$\lim_{R \to \infty} R^2 \int_0^\infty \frac{\sin^2(\pi Rx)}{1 - (Rx)^2} \frac{\tanh(\pi x)}{x} dx = \frac{\pi}{2} \lim_{R \to \infty} \int_0^\infty F(u) H_R(u) du$$
$$= \frac{\pi}{4} \lim_{R \to \infty} \int_0^\infty \frac{H_R(Rx/\pi)}{x} dx$$
$$= \frac{\pi}{4} \lim_{R \to \infty} \int_0^\infty \phi_R(x) dx$$
$$= \frac{\pi}{4} \int_0^\infty \lim_{R \to \infty} \phi_R(x) dx$$
$$= \frac{\pi^2}{2} \int_0^\infty \frac{-h'(x)}{x} dx$$
$$= \frac{\pi^2}{2} \int_0^\infty \frac{\sinh(x) \cosh(x) - x}{x^3 \cosh^2(x)} dx$$

where we used the dominated convergence theorem to pass the limit inside the integral.

To get a rigorous lower bound this last integral, we observe that the integrand is positive so that

$$\frac{\pi^2}{2} \int_0^\infty \frac{\sinh(x)\cosh(x) - x}{x^3\cosh^2(x)} dx \ge \frac{\pi^2}{2} \int_A^B \frac{\sinh(x)\cosh(x) - x}{x^3\cosh^2(x)} dx$$

for any 0 < A < B. For $A = 1/10^6$ and B = 500, interval arithmetic in SageMath certifies that the right-hand side is at least 4.20718596495552 > 4.20718.

9. Multiplicity of the first eigenvalue

9.1. The criterion. We denote the multiplicity of the first positive eigenvalue of the Laplacian on a hyperbolic surface M by $m_1(M)$. The following criterion for bounding m_1 was first stated and proved in [FBP21, Lemma 3.2] in a slightly more general form.

Theorem 9.1. Let M be a closed hyperbolic surface of genus $g \ge 2$ and suppose that f is an admissible function such that

- $f(x) \ge 0$ for all $x \in \mathbb{R}$;
- $\widehat{f}\left(\sqrt{\lambda \frac{1}{4}}\right) \leq 0$ whenever $\lambda \geq \lambda_1(M)$; $\widehat{f}\left(\sqrt{\lambda_1(M) \frac{1}{4}}\right) < 0.$

Then

$$m_1(M) \le \frac{\widehat{f}(i/2) - 2(g-1)\int_0^\infty \widehat{f}(x)x \tanh(\pi x) \, dx}{-\widehat{f}\left(\sqrt{\lambda_1(M) - \frac{1}{4}}\right)}.$$

Proof. Let us write $r_j(M) = \sqrt{\lambda_j(M) - \frac{1}{4}}$. The Selberg trace formula tells us that

$$\begin{aligned} \widehat{f}(i/2) + m_1(M)\widehat{f}(r_1(M)) &\geq \sum_{j=0}^{\infty} \widehat{f}(r_j(M)) \\ &= 2(g-1)\int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(M)} \frac{\Lambda(\gamma)f(\ell(\gamma))}{2\sinh(\ell(\gamma)/2)} \\ &\geq 2(g-1)\int_0^{\infty} \widehat{f}(x)x \tanh(\pi x) \, dx \end{aligned}$$

since $\widehat{f}(r_j(M))$ is non-positive for $j \ge 2$ and f is non-negative on the length spectrum of M. Rearranging yields the desired result.

Observe that since we require f to be non-negative on \mathbb{R} and not constant equal to zero, we automatically have that \hat{f} is positive on the imaginary axis. In particular, Theorem 9.1 cannot be used to prove bounds on the multiplicity of λ_1 for surfaces with $\lambda_1(M) \leq$ 1/4, because then $\sqrt{\lambda_1(M) - 1/4}$ is on the imaginary axis. When applying the linear programming method numerically, it also seems that the resulting bounds tend to infinity as $\lambda_1(M)$ decreases to 1/4 in any fixed genus.

We thus have to use different methods in order to bound $m_1(M)$ when $\lambda_1(M)$ is close to the interval [0, 1/4]. When $\lambda_1(M) \in [0, 1/4]$, a theorem of Otal [Ota08] (later generalized in [OR09]) says that $m_1(M) \leq 2g - 3$. If $\lambda_1(M)$ is a little bit beyond 1/4, then the bound gets slightly worse, namely, we have

$$m_1(M) \le \begin{cases} 2g - 1 & \text{if } \lambda_1(M) \in (1/4, a_g] \\ 2g & \text{if } \lambda_1(M) \in (a_g, b_g) \end{cases}$$

where a_g (resp. b_g) is the smallest eigenvalue of the Laplacian on a hyperbolic disk of area $4\pi(g-1)$ (resp. $2\pi(g-1)$) subject to Dirichlet boundary conditions [FBP21, Theorem 1.1]. Estimates for a_g and b_g are given in [FBP21, Section 2].

We thus use a combination of linear programming bounds and the above inequalities to bound m_1 in a given genus since we need to consider all possible values for λ_1 . Similarly as for kissing numbers, when applying linear programming bounds over an interval I of values for λ_1 , we need to subdivide it into smaller intervals and use a single function f on each subinterval J, taking care to bound $\hat{f}\left(\sqrt{\lambda-\frac{1}{4}}\right)$ for $\lambda \in J$.

9.2. Low genus. The bounds we have obtained on m_1 for g between 2 and 20 are listed in Table 5. They improve upon the previous best upper bound of 2g + 3 from [Sév02] (which applies to all Schrödinger operators on Riemannian surfaces). In genus 2 and 3, our bounds were previously obtained in [FBP21] and we do not repeat these calculations in the ancillary file verify_multiplicity.ipynb that certifies the other values.

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We now discuss lower bounds. In every genus $g \ge 3$, Colbois and Colin de Verdière [CCdV88] constructed closed hyperbolic surfaces satisfying $m_1(M) = \lfloor \frac{1}{2} \left(1 + \sqrt{8g+1} \right) \rfloor$. This has the same order of growth as the maximum of $\lfloor \frac{1}{2} \left(5 + \sqrt{48g+1} \right) \rfloor$ among all closed, connected, orientable Riemannian surfaces of genus g conjectured by Colin de Verdière [CdV87]. We list these conjectured values in the table for comparison. In genus 2, Colin de Verdière's formula comes out to 7 but our upper bound is 6. The conjectured maximum among hyperbolic surfaces is 3.

genus	lower bound	conjecture	LP bound	previous bound
2	3 [FBP21]	3 (hyperbolic), 7 (Riemannian)	6 [FBP21]	7 [Sév02]
3	8 [FBP21]	8	8 [FBP21]	9 [Sév02]
4	4 (§9.2) or 6? [Coo18]	9	10	11 [Sév02]
5	3 [CCdV88]	10	11	13 [Sév02]
6	$4 \left[\text{CCdV88} \right]$	11	13	15 [Sév02]
7	7(9.2)	11	15	17 [Sév02]
8	6(9.2)	12	17	19 [Sév02]
9	4 [CCdV88]	12	19	21 [Sév02]
10	8 (§9.2)	13	20	23 [Sév02]
11	5 [CCdV88]	14	22	25 [Sév02]
12	$5 \left[\text{CCdV88} \right]$	14	24	27 [Sév02]
13	$5 \left[\text{CCdV88} \right]$	15	26	29 [Sév02]
14	12 (§ <mark>9.2</mark>)	15	28	31 [Sév02]
15	7(\$9.2)	15	30	33 [Sév02]
16	8 (§9.2)	16	32	35 [Sév02]
17	6 [CCdV88]	16	34	37 [Sév02]
18	$6 \left[\text{CCdV88} \right]$	17	36	39 [Sév02]
19	8 (§9.2)	17	38	41 [Sév02]
20	6 [CCdV88]	18	40	43 [Sév02]

TABLE 5. Bounds on the maximum of the multiplicity m_1 .

Colbois and Colin de Verdière modelled their examples on graphs and used a transversality argument to control the multiplicity. Another way to obtain lower bounds on multiplicity is to use representation theory [Jen84, BC85, Coo18, FBP21] since the isometry group of a closed hyperbolic surface is finite and acts on the eigenspaces of the Laplacian. This means that if all the irreducible representations of a group have dimension at least d, then the multiplicity of any eigenvalue is at least d. The problem is that there is always the trivial representation of dimension 1, so one must find a way to rule out 1-dimensional real representations from appearing in eigenspaces. Proposition 4.4 in [FBP21] gives such a criterion for kaleidoscopic surfaces, as defined below.

Given integers $2 \leq p \leq q \leq r$, a (p, q, r)-triangle surface (sometimes called quasiplatonic) is a hyperbolic surface of the form \mathbb{H}^2/Γ for some finite-index normal subgroup Γ of the (p, q, r)-triangle group, that is, the group generated by rotations of order p, q, and r around the vertices of a hyperbolic triangle with interior angles $\frac{\pi}{p}$, $\frac{\pi}{q}$, and $\frac{\pi}{r}$ at the corresponding vertices. A kaleidoscopic surface is defined similarly but with the extended triangle group generated by the reflections in the sides of a (p, q, r)-triangle. A hyperbolic surface that admits an orientation-reversing isometry is called *symmetric*, reflexible, or real. Perhaps surprisingly, not every triangle surface is symmetric. In fact, neither of the two Hurwitz surfaces in genus 17 is [Sin74, Theorem 5]. It is therefore desirable to have a criterion for ruling out 1-dimensional real representations for asymmetric surfaces. Even for kaleidoscopic surfaces, the representation theory of Isom⁺ can behave better than that of Isom in some cases. We thus prove the following variant of [FBP21, Proposition 4.4].

Lemma 9.2. If M is a hyperbolic (p, q, r)-triangle surface of area at least $6\pi r \left(\frac{1}{p} - \frac{1}{q} - \frac{1}{r}\right)$, then no 1-dimensional real representations of $\text{Isom}^+(M)$ can occur in the eigenspace corresponding to $\lambda_1(M)$.

Proof. Suppose that f is an eigenfunction contained in a 1-dimensional representation of $\text{Isom}^+(M)$ contained in the λ_1 -eigenspace. Then $\text{Isom}^+(M)$ acts by multiplication by ± 1 on f so the set $f^{-1}(0)$ is invariant. By Courant's nodal domain theorem, the complement of this set (which is a union of analytic curves intersecting transversely [Che76]) has exactly two connected components, so in particular $f^{-1}(0)$ is non-empty. We will show that this leads to a contradiction.

Note that $\operatorname{Isom}^+(M) \geq G := \Gamma/\pi_1(M)$ where Γ is the (p, q, r)-triangle group but a priori the inclusion can be strict (because of inclusions between some triangle groups). We will work G instead of $\operatorname{Isom}^+(M)$ because that makes things simpler. Let T be the any (2, q, r)-triangle used to define Γ , let $\widetilde{\mathcal{T}}$ be the tiling of \mathbb{H}^2 generated by the reflections in the sides of T, and let \mathcal{T} be the projection of $\widetilde{\mathcal{T}}$ to M. Since Γ acts simply transitively on adjacent pairs of triangles in $\widetilde{\mathcal{T}}$ that share a particular kind of side (say joining the vertices of type p and q), so does G on M. In other words, $|G| = |\mathcal{T}|/2$ and

$$area(M) = 2|G| area(T) = 2|G|\pi \left(\frac{1}{p} - \frac{1}{q} - \frac{1}{r}\right)$$

From the hypothesis on area(M), we get $|G| \ge 3r$.

Let Q = M/G, let $\pi : M \to Q$ be the quotient map, and let $F = \pi(f^{-1}(0))$. We have that Q is a sphere with three cone points and F is a finite analytic graph with no isolated points and any vertices of degree 1 are contained in the cone points of Q.

Suppose that a component U of $Q \setminus F$ contains exactly one cone point v of Q. Then $\pi^{-1}(U)$ has |G|/k components, where k is one of p, q, or r, all of which are nodal domains. By the above, we have $|G|/k \ge |G|/r \ge 3$, which contradicts Courant's theorem.

If $Q \setminus F$ has more than two connected components, then f has more than two nodal domains (the preimages of these components). It follows that F contains at most one cycle.

Suppose that F does contain a cycle. Then $Q \setminus F$ has two connected components by the Jordan curve theorem. Since each component contains either 0, 2, or 3 cone points by the above argument, at least one of them, call it U, does not contain any. The quotient map $\pi: M \to Q$ is thus unbranched over U, so that $\pi^{-1}(U)$ has |G| components, all of which are nodal domains. This is again a contradiction since |G| > 2.

We conclude that F is a forest and in particular its complement is connected. Recall that all the leaves of F are contained in the 3 cone points. Moreover, all the cone points must belong to F since it has at least two leaves and if it has only two, then its complement contains exactly one cone point, which is impossible by the above reasoning. We conclude

that F is a tree. Once again, the map π is a covering map over the simply connected domain $Q \setminus F$, so $\pi^{-1}(Q \setminus F)$ has |G| > 2 components, contradiction.

The list of all triangle surfaces in genus 2 to 101 was tabulated by Conder using Magma and is available at [Con15]. We went through all examples in genus 2 to 20, verified if the area hypothesis of [FBP21, Proposition 4.4] or of Lemma 9.2 was satisfied, calculated the character tables for Isom⁺ and Isom using Sage/GAP, and then calculated the second smallest dimension of an irreducible real representation of the group (or rather, a lower bound for it). The cases where the resulting multiplicity is larger than the $\lfloor \frac{1}{2} (1 + \sqrt{8g+1}) \rfloor$ obtained in [CCdV88] are as follows:

- The Bolza surface of genus 2 and type (2, 3, 8) with $m_1 = 3$. Jenni's proof of this fact in [Jen84] contained an error partially corrected in [Coo18] and fixed in [FBP21]. In this case, Isom admits some 2-dimensional irreducible real representations but they can be ruled out with a more careful analysis.
- The Klein quartic of genus 3 and type (2, 3, 7) with $m_1 = 8$ [FBP21] (the representation theory of Isom \cong PGL(2, 7) only gives $m_1 \ge 6$ [Coo18]).
- Bring's curve B of genus 4 and type (2, 4, 5), which satisfies $\text{Isom}^+(B) \cong S_5$ and $\text{Isom}(B) \cong S_5 \times (\mathbb{Z}/2\mathbb{Z})$. While Isom(B) admits some 2-dimensional irreducible real representations, $\text{Isom}^+(B) \cong S_5$ has two real 1-dimensional representations and then irreducible complex representations of dimensions 4, 5, and 6. In particular, its irreducible real representations of dimension more than 1 have real dimension at least 4. By Lemma 9.2, 1-dimensional real representations of $\text{Isom}^+(B) \ge 4$. Numerical evidence suggests that the correct value is $m_1(B) = 6$ [Coo18].
- The Fricke-Macbeath curve F of genus 7 is a Hurwitz surface (by definition, a (2,3,7)-triangle surface) such that $\text{Isom}^+(F) \cong \text{PSL}(2,8) = \text{SL}(2,8)$ (see e.g. [Sin74]). The non-trivial irreducible complex representations of this group have complex dimensions 7, 8, and 9 [Ada02]. It follows from Lemma 9.2 that $m_1(F) \ge 7$. Numerical calculations by Mathieu Pineault suggest that equality holds.
- The (2,3,8)-triangle surface of genus 8 labelled T8.1 in Conder's list satisfies $m_1 \ge 6$ due to the representation theory of Isom⁺.
- The (2, 4, 5)-triangle surface of genus 10 labelled T10.7 in Conder's list satisfies $m_1 \geq 8$ due to the representation theory of Isom.
- In genus 14, there are three distinct Hurwitz surfaces X_j , all with $\text{Isom}^+(X_j) \cong \text{PSL}(2, 13)$ (see [Sin74]) and whenever the group of orientation-preserving isometries of a Hurwitz surface is PSL(2, q), its full group of isometries is PGL(2, q) [BBC⁺96, Remark 2.3(2)]. The group $\text{PGL}(2, 13) \cong \text{Isom}(X_j)$ has two 1-dimensional real representations (the trivial one and the sign representation) and then irreducible complex representations of dimensions 12, 13, and 14 [Ada02]. By [FBP21, Proposition 4.4], the three surfaces X_j all satisfy $m_1 \geq 12$.
- The (2,3,9)-triangle surface of genus 15 labelled T15.1 in Conder's list satisfies $m_1 \ge 7$ due to the representation theory of Isom⁺.
- The (2,3,8)-triangle surface of genus 16 labelled T16.1 in Conder's list satisfies $m_1 \ge 8$ due to the representation theory of Isom⁺.
- The (2, 4, 5)-triangle surface of genus 19 labelled T19.3 in Conder's list satisfies $m_1 \ge 8$ due to the representation theory of Isom.

The ancillary file verify_multiplicity_examples.ipynb contains the computer code that found these examples. Interestingly, in most genera the bound of Colbois and Colin de Verdière is matched by some triangle surface. However, since the sum of the squares of the degrees of the irreducible representations of a group is equal to the order of the group, and since the isometry group of a closed hyperbolic surface of genus g has order at most 168(g-1) [Hur92], the best this method could give is still $O(\sqrt{g})$.

Back to upper bounds, note that the linear programming bounds listed in Table 5 never go below 2g in the range considered here, so we can assume that λ_1 is between b_g and the upper bounds from Table 3 to prove them. For genus 6 onwards, our bounds are worse when λ_1 is close to (the estimate for) b_g than at the upper bound on λ_1 . For example in genus 20, we obtain an upper bound of 23 instead of 40 when $\lambda_1 \in [0.750384, 0.91139]$. Our intuition is that m_1 should be maximized among local maximizers of λ_1 . The fact that the bound on m_1 increases when λ_1 decreases is an artefact of the method and is not necessarily representative of the reality.

9.3. Asymptotics. In higher genus, we will decrease Sévennec's upper bound by 4. We start by proving a sublinear bound on m_1 under the assumption that λ_1 is fairly large. In the statement below, j_{α} stands for the first positive zero of the Bessel function J_{α} . The k-th positive zero is denoted $j_{\alpha,k}$.

Proposition 9.3. For every $p_1, p_2 \in (j_0, \pi]$, there exists a constant C and some $g_0 \ge 2$ such that if $g \ge g_0$ and M is a closed hyperbolic surface of genus g with

$$\lambda_1(M) \in \left\lfloor \frac{1}{4} + \left(\frac{p_1}{\log(g)}\right)^2, \frac{1}{4} + \left(\frac{p_2}{\log(g)}\right)^2 \right\rfloor$$

then $m_1(M) \leq Cg/\log(g)^3$.

Proof. We start by picking two constants α and c that only depend on p_1 and p_2 (and not on g).

For every $k \in \mathbb{N}$, the k-th zero $j_{\alpha,k}$ of J_{α} is increasing as a function of α provided that $\alpha \geq 0$ [Wat95, p.507]. Together with the interlacing property of the Bessel zeros [Wat95, p.479], this implies that

$$j_{0,2} > j_{1,1} > j_{1/2,1} = \pi > j_{0,1} = j_0$$

In fact, much better numerics are known, namely, $j_0 \approx 2.4048$ and $j_{0,2} \approx 5.5201$ [DK27]. In particular, we have

$$\frac{p_1}{p_2} > \frac{j_0}{\pi} > \frac{j_0}{j_{0,2}}$$

Since $j_{\alpha,k}$ depends continuously on α for every k [Wat95, p.507], there exists some α in (0, 1/2) such that

$$p_1 > j_{\alpha}$$
 and $\frac{p_1}{p_2} > \frac{j_{\alpha}}{j_{\alpha,2}}$

We then take any $c \in (j_{\alpha}/p_1, \min\{j_{\alpha,2}/p_2, 1\})$. This interval is not empty since $j_{\alpha}/p_1 < 1$ and $j_{\alpha}/p_1 < j_{\alpha,2}/p_2$.

We set $R = c \log(g)$ and apply Theorem 9.1 with the function f such that

$$\widehat{f}(x) = \psi_{\alpha}(Rx) = \frac{J_{\alpha}(Rx)^2}{(Rx)^{2\alpha}(1 - R^2x^2/j_{\alpha}^2)}$$

The function \hat{f} is positive-definite by [GIT20, Remark 1.1] and

$$\mu := \min\left\{-\widehat{f}(x) : x \in \left[\frac{p_1}{\log(g)}, \frac{p_2}{\log(g)}\right]\right\} = \min\left\{-\psi_{\alpha}(y) : y \in [c\,p_1, c\,p_2]\right\}$$

is a positive constant that only depends on our choice of parameters, but not on the genus. Indeed, ψ_{α} is continuous and strictly negative on $(j_{\alpha}, j_{\alpha,2}) \supset [c p_1, c p_2]$.

By the asymptotic behaviour of Bessel functions along the imaginary axis (5.1), we have

$$\widehat{f}(i/2) \sim \frac{1}{\pi} \frac{e^R}{R^{2\alpha+1}(1+R^2/(4j_{\alpha}^2))} = o\left(\frac{g}{\log(g)^3}\right)$$

as $g \to \infty$.

if $\sqrt{}$

For the integral term, we make a change of variable to find that

$$\left| \int_0^\infty \widehat{f}(x)x \tanh(\pi x) \, dx \right| = \left| \int_0^\infty \psi_\alpha(Rx)x \tanh(\pi x) \, dx \right|$$
$$= \left| \frac{\pi}{R^3} \int_0^\infty \psi_\alpha(y)y^2 \frac{\tanh(\pi y/R)}{\pi y/R} \, dy \right|$$
$$\leq \frac{\pi}{R^3} \int_0^\infty |\psi_\alpha(y)|y^2 \, dy$$
$$= O\left(\frac{1}{\log(g)^3}\right).$$

Indeed, $|\psi_{\alpha}(y)|y^2$ is integrable by the asymptotic estimate (5.2) since $\alpha > 0$.

By Theorem 9.1 applied with the function f, we obtain

$$m_1(M) \le \frac{1}{\mu} \left(\widehat{f}(i/2) - 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) dx \right) = O\left(\frac{g}{\log(g)^3}\right)$$
$$\overline{\lambda_1(M) - 1/4} \in \left[\frac{p_1}{\log(g)}, \frac{p_2}{\log(g)}\right] \text{ and if } g \text{ is sufficiently large.} \qquad \square$$

Remark 9.4. It is easy to modify the proof of Theorem 9.1 to bound the total number of eigenvalues in an interval [a, b] assuming that $\lambda_1(M) \in [a, b]$ (see [FBP21, Lemma 3.2]). This more general version of the criterion implies that under the same hypotheses as above, the total number of eigenvalues contained in the given interval is bounded by the same quantity $Cg/\log(g)^3$.

See [GLMST21, Corollary 1.7] for a similar result showing that for any compact interval I contained in $[0, 1/4) \cup (1/4, \infty)$, the multiplicity of any eigenvalue in I grows at most sublinearly with the genus g with probability tending to 1 as $g \to \infty$ with respect to the Weil–Petersson measure. See also [Mon22, Corollary 6].

We then combine the sublinear upper bound from Proposition 9.3 with a previous (linear) upper bound from [FBP21] for slightly smaller λ_1 to obtain a global upper bound on m_1 in large genus.

Theorem 9.5. There exists some $g_0 \ge 2$ such that every closed hyperbolic surface M of genus $g \ge g_0$ satisfies

$$m_1(M) \le 2g - 1.$$

Proof. By Theorem 8.3, we can assume that

$$\lambda_1(M) \le \frac{1}{4} + \left(\frac{\pi}{\log(g) + 0.7436}\right)^2 \le \frac{1}{4} + \left(\frac{\pi}{\log(g)}\right)^2.$$

Now pick any $p \in (j_0, \pi)$. Proposition 9.3 implies (a stronger version of) the result if

$$\lambda_1(M) \ge \frac{1}{4} + \left(\frac{p}{\log(g)}\right)^2$$

It remains to consider the case where $\lambda_1(M)$ is smaller than this bound. This case was already handled in [FBP21, Theorem 1.1], which states that $m_1(M) \leq 2g - 1$ whenever $\lambda_1(M) \leq a_g$, where a_g is the smallest Dirichlet eigenvalue on a hyperbolic disk of area $4\pi(g-1)$ and hence of radius $2 \operatorname{arcsinh}(\sqrt{g-1})$. From Savo's inequality (8.3), we have

$$a_g - \frac{1}{4} \geq \frac{\pi^2}{4\operatorname{arcsinh}(\sqrt{g-1})^2} - \frac{\pi^2}{2\operatorname{arcsinh}(\sqrt{g-1})^3} \sim \frac{\pi^2}{\log(g)^2} \quad \text{as } g \to \infty.$$

In particular, when q is large enough we have

$$a_g > \frac{1}{4} + \left(\frac{p}{\log(g)}\right)^2.$$

As such, all possibilities for $\lambda_1(M)$ are covered and the inequality is proved.

Note that the proof relied partly on [FBP21, Theorem 1.1] whose proof is based on the results of Sévennec [Sév02]. Different methods would be required to prove sublinear upper bounds.

10. Small eigenvalues

10.1. The criterion. We start by proving a general criterion for bounding the number $N_{[0,L]}(M)$ of eigenvalues of Δ_M in the interval [0, L] for any L > 0.

Theorem 10.1. Let L, R > 0 and suppose that $f : \mathbb{R} \to \mathbb{R}$ is an admissible function such that

- f(x) ≤ 0 for every x ≥ R;
 f̂(ξ) ≥ 0 for every ξ ∈ ℝ;
 f̂(√λ − 1/4) ≥ 1 for every λ ∈ [0, L];

Then every closed hyperbolic surface M of genus $g \ge 2$ with $sys(M) \ge R$ satisfies

$$N_{[0,L]}(M) \le 2(g-1) \int_0^\infty \widehat{f}(x) x \tanh(\pi x) \, dx + 1 - \widehat{f}(i/2).$$

Proof. By hypothesis, $\widehat{f}(r_j(M))$ is at least 1 for all eigenvalues $\lambda_j(M)$ in the interval [0, L]and non-negative at all eigenvalues. Also note that $\lambda_0(M) = 0 \in [0, L]$ corresponds to the

term $\widehat{f}(i/2)$. We thus have

$$\begin{split} N_{[0,L]}(M) - 1 + \widehat{f}(i/2) &\leq \sum_{j=0}^{\infty} \widehat{f}(r_j(M)) \\ &= 2(g-1) \int_0^{\infty} \widehat{f}(x) x \tanh(\pi x) \, dx + \frac{1}{\sqrt{2\pi}} \sum_{\gamma \in \mathcal{C}(M)} \frac{\Lambda(\gamma) f(\ell(\gamma))}{2 \sinh(\ell(\gamma)/2)} \\ &\leq 2(g-1) \int_0^{\infty} \widehat{f}(x) x \tanh(\pi x) \, dx \end{split}$$

by the Selberg trace formula, where the last inequality is because the geometric terms are non-positive by hypothesis. $\hfill \Box$

Recall that an eigenvalue of the Laplacian on M is *small* if it belongs to the interval [0, 1/4]. The number $N_{\text{small}}(M)$ of small eigenvalues of M is therefore equal to $N_{[0,1/4]}(M)$. Also note that $\lambda \in [0, 1/4]$ if and only if $\sqrt{\lambda - 1/4} \in i[-1/2, 1/2]$. In practice, we will often use the weaker bound

$$\frac{N_{\text{small}}(M)}{2(g-1)} \le \inf_{f} \int_{0}^{\infty} \widehat{f}(x) x \tanh(\pi x) \, dx$$

(where the infimum is over the functions f that satisfy the hypotheses of the theorem) instead of the one given in Theorem 10.1 because the right-hand side only depends on sys(M) and not on the genus. A theorem of Otal and Rosas [OR09] states that the left-hand side is always bounded above by 1, and this sharp for surfaces with a very short pants decomposition [Bus77] (see also [Bus10, Theorem 8.1.3]). Our goal is therefore to find the smallest value of R = sys(M) for which the right-hand side becomes smaller than 1 and to estimate how it decreases as the systole increases.

10.2. Asymptotics. We start with the asymptotics instead of the numerics for N_{small} because they are easier to obtain than for the other invariants and they give us a point of comparison for the numerics.

Theorem 10.2. If M is a closed hyperbolic surface of genus $g \ge 2$, then

$$N_{\text{small}}(M) < \min\left(\frac{24\pi^2(g-1)}{\text{sys}(M)^3}, \frac{16(g-1)}{\text{sys}(M)^2}\right)$$

Proof. Let $f : \mathbb{R} \to \mathbb{R}$ be an even, non-negative, admissible (hence continuous), positivedefinite function supported in [-1, 1] normalized so that $\hat{f}(0) = 1$.

We then use the function $f_R(x) = f(x/R)/R$ in Theorem 10.1 with L = 1/4 and R = sys(M). The first and second bullet points in the statement of the theorem are satisfied by hypothesis on f while the third one follows from the non-negativity of f (and hence of f_R). Indeed,

$$\widehat{f_R}(it) = \sqrt{\frac{2}{\pi}} \int_0^\infty f_R(x) \cosh(tx) \, dx \ge \sqrt{\frac{2}{\pi}} \int_0^\infty f_R(x) \, dx = \widehat{f_R}(0) = \widehat{f}(0) = 1$$

for every $t \in \mathbb{R}$, with equality only if t = 0.

If \widehat{f} has a finite first moment on $[0,\infty)$, then we can estimate

$$\int_0^\infty \widehat{f_R}(x)x \tanh(\pi x) \, dx = \int_0^\infty \widehat{f}(Rx)x \tanh(\pi x) \, dx$$
$$< \int_0^\infty \widehat{f}(Rx)x \, dx = \frac{1}{R^2} \int_0^\infty \widehat{f}(y)y \, dy$$

The function f determined by $\hat{f}(x) = \eta_{\alpha}(x/2)^2$, where

$$\eta_{\alpha}(x) = 2^{\alpha} \Gamma(\alpha + 1) \frac{J_{\alpha}(x)}{x^{\alpha}}$$

is the normalized Bessel function, satisfies all the necessary requirements provided that $\alpha > 1/2$ (see Section 5). We then compute

$$\int_0^\infty \widehat{f}(y)y \, dy = \int_0^\infty \eta_\alpha (y/2)^2 y \, dy = 4 \int_0^\infty \eta_\alpha (x)^2 x \, dx.$$

The recurrence formulae for Bessel functions [Wat95, p.2] imply that

$$\eta'_{\alpha}(x) = \frac{-x}{2(\alpha+1)}\eta_{\alpha+1}(x) = \frac{2\alpha}{x}(\eta_{\alpha-1}(x) - \eta_{\alpha}(x))$$

and from this it is easy to check that

$$\frac{1}{2-4\alpha} \left(4\alpha^2 \eta_{\alpha-1}(x)^2 + \eta_{\alpha}(x)^2 x^2 \right)$$

is a primitive of $\eta_{\alpha}(x)^2 x$ which vanishes at infinity. We thus have

$$\int_{0}^{\infty} \widehat{f}(y) y \, dy = 4 \int_{0}^{\infty} \eta_{\alpha}(x)^{2} x \, dx = \frac{16\alpha^{2}}{4\alpha - 2}$$

since $\eta_{\beta}(0) = 1$ for every β . It is easy to check that this quantity is at least 8 with equality only if $\alpha = 1$. For that parameter, the resulting inequality is

$$N_{\rm small}(M) < \frac{16(g-1)}{\rm sys}(M)^2$$

if we ignore the term $1 - \widehat{f_R}(i/2) < 0$.

We can also write

$$\int_0^\infty \widehat{f_R}(x)x \tanh(\pi x) \, dx = \int_0^\infty \widehat{f}(Rx)x \tanh(\pi x) \, dx$$
$$< \pi \int_0^\infty \widehat{f}(Rx)x^2 \, dx$$
$$= \frac{\pi}{R^3} \int_0^\infty \widehat{f}(y)y^2 \, dy.$$

provided that \widehat{f} has a finite second moment.

With the same function $\widehat{f}(x) = \eta_{\alpha}(x/2)^2$ as before (but with $\alpha > 1$ this time), we need to compute

$$\int_0^\infty \widehat{f}(y)y^2 \, dy = \int_0^\infty \eta_\alpha (y/2)^2 y^2 \, dy = 8 \int_0^\infty \eta_\alpha (x)^2 x^2 \, dx.$$

Integration by parts with u = x and $dv = \eta_{\alpha}(x)^2 x \, dx$ yields

$$\int_0^\infty \eta_\alpha(x)^2 x^2 \, dx = \frac{1}{4\alpha - 2} \int_0^\infty \left(4\alpha^2 \eta_{\alpha - 1}(x)^2 + \eta_\alpha(x)^2 x^2 \right) \, dx$$

and hence

$$\int_0^\infty \eta_{\alpha}(x)^2 x^2 \, dx = \frac{4\alpha^2}{4\alpha - 3} \int_0^\infty \eta_{\alpha - 1}(x)^2 dx.$$

Recall that

$$\widehat{\eta_{\beta}}(t) = \frac{\sqrt{\pi/2}}{B(\frac{1}{2}, \alpha + \frac{1}{2})} \operatorname{rect}(t/2)(1 - t^2)^{\beta - 1/2}$$

for every $\beta > 1/2$ and every $t \in \mathbb{R}$, where rect is the characteristic function of [-1/2, 1/2]and B is the Beta function. By the convolution formula, we have

$$\begin{split} \int_{0}^{\infty} \eta_{\beta}(x)^{2} dx &= \sqrt{\frac{\pi}{2}} \, \widehat{\eta}_{\beta}^{2}(0) = \frac{1}{2} \, \widehat{\eta_{\beta}} * \widehat{\eta_{\beta}}(0) \\ &= \frac{1}{2} \left(\frac{\sqrt{\pi/2}}{B(\frac{1}{2}, \beta + \frac{1}{2})} \right)^{2} \int_{-1}^{1} (1 - t^{2})^{2\beta - 1} dt \\ &= \frac{\pi/2}{B(\frac{1}{2}, \beta + \frac{1}{2})^{2}} \int_{0}^{1} (1 - t^{2})^{2\beta - 1} dt \\ &= \frac{\pi/4}{B(\frac{1}{2}, \beta + \frac{1}{2})^{2}} \int_{0}^{1} u^{1/2 - 1} (1 - u)^{2\beta - 1} du \\ &= \frac{\pi/4}{B(\frac{1}{2}, \beta + \frac{1}{2})^{2}} B(\frac{1}{2}, 2\beta) \\ &= \frac{\pi}{4} \left(\frac{\Gamma(\beta + 1)}{\Gamma(1/2)\Gamma(\beta + 1/2)} \right)^{2} \frac{\Gamma(1/2)\Gamma(2\beta)}{\Gamma(2\beta + 1/2)} \end{split}$$

Using the recursion $\Gamma(z+1) = z\Gamma(z)$, Legendre's duplication formula

$$\Gamma(z)\Gamma(z+1/2) = 2^{1-2z}\sqrt{\pi}\Gamma(2z),$$

and the special value $\Gamma(1/2) = \sqrt{\pi}$, the above simplifies to

$$\int_0^\infty \eta_\beta(x)^2 dx = \frac{2^{8\beta-3}}{\pi} \frac{\beta^2 \Gamma(\beta)^4}{\Gamma(4\beta)}$$

Returning to the original problem, we have

$$\int_0^\infty \widehat{f}(y)y^2 \, dy = 8 \int_0^\infty \eta_\alpha(x)^2 x^2 \, dx = \frac{32\alpha^2}{4\alpha - 3} \int_0^\infty \eta_{\alpha - 1}(x)^2 dx$$
$$= \frac{2^{8\alpha - 6}}{\pi} \frac{\alpha^2(\alpha - 1)^2}{4\alpha - 3} \frac{\Gamma(\alpha - 1)^4}{\Gamma(4\alpha - 4)}.$$

One can check that this function is minimized at $\alpha = 3/2$, where it takes the value 12π . The resulting inequality is

$$N_{\text{small}}(M) < \frac{24\pi^2(g-1)}{\text{sys}(M)^3}.$$

Remark 10.3. Note that $24\pi^2/x^3 < 16/x^2$ if and only if $x > 3\pi^2/2 \approx 14.8044$ (for x > 0).

We now compare our inequality

$$N_{\text{small}}(M) < \frac{24\pi^2(g-1)}{\text{sys}(M)^3}$$

with Huber's inequality

$$N_{\text{small}}(M) \le \frac{3\pi^2(g-1)}{8(\log(\cosh(\text{sys}(M)/4)))^3}$$

from [Hub76]. Since $\cosh(x) < e^x$ for every x > 0, we have $\log(\cosh x) < x$ and hence

$$\frac{3\pi^2(g-1)}{8(\log(\cosh(\operatorname{sys}(M)/4)))^3} > \frac{3\pi^2(g-1)}{8\operatorname{sys}(M)^3/64} = \frac{24\pi^2(g-1)}{\operatorname{sys}(M)^3},$$

so that our bound is better but only slightly. Indeed, the inequality $\cosh(x) > e^x/2$ implies that the factors that multiply (g-1) in the two inequalities are asymptotic to each other as $\operatorname{sys}(M) \to \infty$ despite the fact that the two proofs use different functions (Huber uses Legendre functions while we use Bessel functions).

From our inequality

$$N_{\text{small}}(M) < \frac{16(g-1)}{\text{sys}(M)^2}$$

it follows that if $sys(M) \ge \sqrt{8}$, then $N_{small}(M) < 2g - 2$. In [Hub76], Huber also proves the inequality

$$N_{\text{small}}(M) \le \frac{g-1}{2\log(\cosh(\text{sys}(M)/4))},$$

which implies that $N_{\text{small}}(M) < 2g - 2$ as soon as

$$sys(M) > 4 \operatorname{arccosh}(e^{1/4}) \approx 2.947618.$$

This is better than the constant 3.46 recently obtained in [Jam21], but not as good as $\sqrt{8} \approx 2.828427$. In fact, one can show that

$$\frac{16}{x^2} < \frac{1}{2\log(\cosh(x/4))}$$

for every x > 0, which means that our bound is better than Huber's for every value of sys(M). We will further decrease the lower bound on the systole sufficient to improve upon the inequality $N_{small}(M) \leq 2g - 2$ of Otal and Rosas in the next subsection.

10.3. Numerical results for small systole. Unsurprisingly, numerical optimization yields better results than Theorem 10.2 when the systole is relatively small. For example, the resulting bounds show that $N_{\text{small}}(M) < 2g - 2$ as soon as $\text{sys}(M) \ge 2.317$ and that $N_{\text{small}}(M) < g - 1$ if $\text{sys}(M) \ge 3.234$. A list of lower bounds on sys(M) and the upper bounds they imply on $N_{\text{small}}(M)/(2g - 2)$ is given in Table 6. The verification of these values is done in the ancillary file verify_nsmall.ipynb. To produce the plot in Figure 2a, we used these values as well as the bounds produced at many other points and took a spline through this list of points. Thus, the plot itself is not rigorous, but the table is.

lower bound on $sys(M)$	strict upper bound on $N_{\text{small}}(M)/(2g-2)$
2.317	1
3.234	1/2
3.919	1/3
4.486	1/4
4.978	1/5
5.409	1/6
5.818	1/7
6.180	1/8
6.505	1/9
6.894	1/10

TABLE 6. Lower bounds on sys(M) sufficient for $N_{small}(M)/(2g-2)$ to be strictly smaller than given values.

10.4. Ramanujan surfaces. Borrowing terminology from graph theory, we say that a hyperbolic surface M of finite area is Ramanujan if $\lambda_1(M) \ge 1/4$. We will also say that M is strictly Ramanujan if $\lambda_1(M) > 1/4$. Selberg's eigenvalue conjecture [Sel65] states that all congruence covers of the modular curve are Ramanujan. A related question is whether there exist closed Ramanujan surfaces in every genus (see [Mon15, Question 1.1] and [Wri20, Problem 10.4]). Thanks to the work of Hide and Magee [HM21], we now know that in large genus, there exist closed surfaces that are nearly Ramanujan in the sense that their first eigenvalue is arbitrarily close to 1/4.

Observe that

$$\lambda_1(M) > L \iff N_{[0,L]}(M) = 1 \iff N_{[0,L]}(M) < 2.$$

This means that one can prove lower bounds on λ_1 by bounding $N_{[0,L]}$ from above. In particular,

$$\lambda_1(M) > \frac{1}{4} \iff N_{\text{small}}(M) < 2 \iff \frac{N_{\text{small}}(M)}{2g-2} < \frac{1}{g-1}$$

so the values in Table 6 give lower bounds on the systole that are sufficient for surfaces to be strictly Ramanujan in genus 2 to 11. However, in obtaining these values we discarded the term $1 - \hat{f}(i/2)$ appearing in Theorem 10.1. The lower bounds on the systole (sufficient to be strictly Ramanujan) that we have obtained by taking this term into account are listed in Table 7 for g between 2 and 20 and plotted in Figure 2b. The corresponding ancillary file is verify_ramanujan.ipynb. According to Table 1, there exist hyperbolic surfaces with systole larger than these bounds, hence strictly Ramanujan, in genus 2 to 7, 14, and 17. For these specific surfaces, we can increase L further as long as $N_{[0,L]} < 2$ to obtain improved lower bounds on λ_1 still based only on the systole. The resulting bounds are listed in Table 3 except in genus 2 to 4 where better data was already available. The corresponding ancillary file is verify_ramanujan_examples.ipynb. These bounds are rigorous modulo proving that the lower bounds on the systole in Table 1 are correct (as pointed out earlier, some are rigorous but not all).

It seems very likely that surfaces with systole larger than the values listed in Table 7 exist in the remaining genera up to 20 as well. However, our numerical experiments suggest that this method cannot prove that the next Hurwitz surface (of genus 118) is Ramanujan

genus	lower bound on sys(M) sufficient for $\lambda_1(M) > 1/4$	record systole
2	2.315	3.057141
3	3.218	3.983304
4	3.867	4.624499
5	4.380	4.91456
6	4.803	5.109
7	5.168	5.796298
8	5.482	
9	5.760	
10	6.010	
11	6.236	
12	6.443	
13	6.632	
14	6.808	6.887905
15	6.971	
16	7.124	
17	7.268	7.609407
18	7.403	
19	7.531	
20	7.651	

TABLE 7. Lower bounds on sys(M) sufficient for $\lambda_1(M) > 1/4$.

using only its systole. Based on Figure 1a, Figure 2b, and Theorem 6.4, it seems reasonable to make the following conjecture.

Conjecture 10.4. There exist constants $c_1 > c_2$ such that the smallest upper bound on the systole that can be obtained from Theorem 6.1 and the smallest lower bound on the systole sufficient to prove $N_{\text{small}} < 2$ (equivalently $\lambda_1 > 1/4$) using Theorem 10.1 are of the form

 $2\log(g) + c_i + o(1)$ as $g \to \infty$.

It is unknown if there exist closed hyperbolic surfaces with systole asymptotic to $r \log(g)$ for any r > 4/3, so even if the conjecture is true it is unlikely to be good enough to prove the existence of Ramanujan surfaces in large genus.

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